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PART III. RELATIVIST	IC GAS DINAMICS	
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1. Basic thermodynamic relationshi	ps,	351
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A UTHOR :	Kaplan, S.A.	507/115-58 - 1-2/50
TITLE:		ng Laboratories for Linear Measure- zmeritel'nykh laboratoriy dlya
PERIODICAL:	Izmeritel'naya tekhnika, 19	958, Nr 1, pp 5-7 (USSR)
ABSTRACT :	ard layout for measurement tion machine-building plan of Measures and Measuring ly the Ministry of Constru- The information includes a list of required laborator; plan for this equipment wi be used for estimating the	led information on the basic stand- laboratories of non-mass produc- ts, worked out by the Department Devices of Vptistroydormash (former- ction and Road Machine-Building). building layout, a basic minimum y equipment, and a distribution thin the building. Equations to required quantities of instruments are given. There are 2 tables and
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sov/169-59-4-4068 Translation from: Referativnyy zhurnal, Geofizika, 1959, Nr 4, p 127 (USSR) Kaplan, S.A. 11 Magnetic Gas Dynamics and the Problems of Cosmogony AUTHOR 1 V sb.: Vopr. kosmogonii. Vol 6, Moscow, AS USSR, 1958, pp 238-TITLE: PERIODICAL 264 (Engl. Res.) The basic principles and the general results obtained in the magnetic gas dynamics are discussed: 1) the existence of the Tadhesion" integral; 2) the description of the "entanglement" ABSTRACT: and the "disentanglement" of the magnetic force lines; 3) the increase of the density of magnetic energy in gas dynamic shock waves; 4) the notion of the gas magnetic turbulence. Certain cosmogonic hypotheses are discussed, in which the methods and results of magnetic gas dynamics are used to some degree: a) the hypothesis of the connection between spiral arms and the regular magnetic field; b) the hypothesis of the formation of interstellar gas clouds as individual vortices - nuclei of the Vinterstellar magnetic turbulence; c) the hypothesis of the Card 1/2

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KAPLAN, S.A. MENA EREALS Approximate calculation of ephemerides and the determination of orbits of artificial earth satellites. Astron.tsir. no.192:5-8 (MIRA 11:10) Ny 158. 1.L'vovskaya stantsiya nablyudeniy Iskusstvennykh Sputnikov Zemli. (Artificial satellites)



FALSE AMERINATERITY, David Al'bertovich; KAFLAN, S.A., retsensent; SAMEONEMEO, L.V., red.; MERMAKOVA, Ye.A., teknn.red.
[Internal physical processes in stars] Fisicheskie protsessy ymutri svesd. Moskva, Gos.isd-vo fisiko-matem.lit-ry, 1959. (MIRA 13:3)
(Astrophysics)

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KAPIAN, S.A.; LOGVINEHKO, A.A. [Logvynenko, O.O.]; PODSTRIGACH, T.S. [Pidstryhach, T.S.] ļ Calculation of gasomagnetic shock wave parameters. Ukr.fis. shur. 4 no.4:438-442 J1-Ag '59. (HIRA 1) (HIRA 13:4) 1. L'vovskiy gosudarstvennyy universitet im. Iv. Franko. (Shock waves) 14

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SOV/33-36-2-21/27 Son the Correlation Between the Observed Differences of the Degree of Interstellar Polarization and the Angular Distance of the Corresponding Points on the Celestial Sphere There are 2 references, 1 of which is Soviet, and 1 American. ASSOCIATION: L'vovskaya astronomicheskaya observatoriya (L'vov Astronomical Observatory) SUBMITTED: June 2, 1958 Card 2/2

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3(1),10(1) AUTHORS:	Kaplan, S.A., and Klimishin, I.A. SOV/33-36-3-3/29
TITLE:	Shock Waves in Stellar Envelopes
PERIODICAL	Astronomicheskiy zhurnal, 1959, Vol 36, Nr 3, pp 410-421 (USSR)
ABSTRACT:	The authors consider physical properties of stellar shock waves, the possibility of separation of the envelopes etc. The shock waves are assumed to by stationary, at the other hand, the inter- action with the radiation is considered. §1 contains the derivation of the formula for the Hugonict-adiabatic curves and other general relations. Because of the complicatedness of the obtained system in the following paragraphs, the authors restrict themselves to especially interesting special cases. §2 is devoted to the so-called detonation-recombination shock waves in a gas- radiation-mixture (these waves move due to the energy liberated during the recombination of ions in the wave front). The waves are described by the equations
	$x^{4}(6\Gamma_{2}-3\Gamma_{2}\beta_{2}-1) - x^{2}[3\Gamma_{2}^{2}(2-\beta_{1}) + 6\Gamma_{2}(\beta_{2}-\beta_{1}) + 8-3\beta_{1} +$
Card 1/3	$+ \frac{2\mathbf{q}\mathbf{q}_{1}}{P_{1}} (\Gamma_{2}+1)^{2} + \Gamma_{2}(\Gamma_{2}+8-3\beta_{2}) = 0$

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Shock Waves in Stellar Envelopes

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$$\frac{P_2}{P_1} = \frac{x^2 + 1}{\Gamma_2 + 1} ; \quad \frac{9_2}{9_1} = \frac{(\Gamma_2 + 1)x^2}{\Gamma_2(x^2 + 1)} , \quad x = \frac{v_1}{\sqrt{P_1/9_1}}$$
$$\Gamma = \beta + \frac{4(4 - 3\beta)^2}{3\beta + 24(1 - \beta)} ; \quad \delta = \frac{5}{3} .$$

The indices 1 and 2 denote the values before and after the passage of the wave. $\beta = {}^{P}g/P$ is the ratio of the gas pressure to the full pressure; q is the set of nascent energy, v is the gas velocity with respect to the front of the wave, g is the density. The system is solved by successive approximation, where the fact, that the detonation-recombination waves are weak, facilitates the solution. In §3 the conditions are found under which a separation of the outer part of the envelope of a red giant taking place with a small velocity is possible. An undisturbed separation of an envelope mass amounting ca. $10^{-3} \div 10^{-5}$ solar masses is possible e.g. if the radius of the giant is 80 - 100 times greater than the solar radius, the mass of the giant nearly equals the solar mass and its absolute magnitude is $-4^{m} \cdot 5$ or

Card 2/3 .5^m.8. The velocity of the separating part is 50 km/sec, the

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Shock Waves in Stellar Envelopes SOV/33-36-3-3/29 velocity of the shock wave 110 km/sec. The place of the separation lies nearly in the center of the radius. §4 treats the influence of radiative cooling on the parameters of a shock wave. It is stated that this influence is essential even at optical depths of \sim 30 and that it leads to a 10 - 100-fold diminution of the temperature behind the wave. §5 is devoted to the properties of shock waves in a degenerated gas. There are 15 references, 12 of which are Soviet, 1 American, 1 English, and 1 German. ASSOCIATION: L'vovskaya astronomicheskaya observatoriya (L'vov Astronomical Observatory) SUBMITTED: June 2, 1958 Card 3/3

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V.CELEBS

AUTHOR:	Kaplan, S. A.	5 0¥/56-36-6-44/66
TITLE:	On the "Larmoron" Th plasmy)	eory of the Plasma (O "larmoronnoy" teorii
PERIODICAL:	Zhurnal eksperimenta Nr 6, pp 1927 -	l'noy i teoreticheskoy fiziki, 1959, Vol 36 1928 (USSR)
ABSTRACT :	particles into the p "larmorons"; they ar ment µ, which are in	of introducing quasiparticles into the mod- ave rise to the introduction of effective lasma theory. In this case they are called e effective particles with the magnetic mo- the leading center of the Larmor motion of $mv_{\perp}^2/2H$ (m is the particle mass and v_{\perp} - the
Car d 1/2	the larmoron is equal Spitser, Belyayev et of larmorons, without author of the present equations, which repr motion velocity of 14	hich is perpendicular to H). The energy of to the total energy of the real particles al. (Refs 1,2) already used the conception t, however, defining it so rigorously. The t "Letter to the Editor" gives a number of resent the components of the progressive armorons (the components u, v, w, where w a, and u and v are perpendicular to it) as

	functions of the velocity compon the larmor frequency and larmoro expressions for the total energy city distribution, the equation The results obtained are discuss Glauberman for discussing the su erences.	n life time, and further also of the larmorons, the velo- of motion and its solution.
ASSOCIATION:	L'vovskiy gosudarstvennyy univer	sitet (L'vov State University)
SUBMITTED:	January 30, 1959	
Card 2/2		

Tees-

KAPLAN, S. A. (L'vov)

"The Structure of a Shock Wave in Plasma."

report presented at the First All-Union Congress on Theoretical and Applied Mechanics, Moscow, 27 Jan - 3 Feb 1960.

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KAPLAH, S.A.; KLIMOVSKAYA, A.I.

Equation of the motion of an artificial earth satellite in horisontal coordinates. Biul.sta.opt.nabl.isk.sput.Zem. no.1: 10-12 '60. (MIRA 13:5)

1. L'vovskaya stantsiya nablyudeniy iskusstveniykh sputnikov Zemli. (Artificial satellites)

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	SOV/33-37-1-2/31
AUTHORS:	Kaplan, S. A., Klimishin, I. A., Sivers, V. N.
TITLE:	A Theory of Light Scattering In a Medlum With a Moving Boundary
PERIODICAL:	Astronomicheskiy zhurnal, 1950, Vol 37, Nr 1, pp 9-15 (USSR)
ABSTRACT:	When the motion of a gas under cosmical conditions is considered, it is frequently necessary to take into account its interaction with radiation. Usually, the problem is studied by combining the equations of motion with the equations of radiative transfer; moreover, only the case of a steady boundary is considered, while actually the scattering occurs either before or after the light quantum passes through a moving boundary. Consequently, before any modern theory of light scatter- ing is applied to hydrodynamic problems it is necessary to develop a theory of scattering in a medium with moving boundaries. This is the problem of the present
Card 1/4	authors. The following notations are used: k, the

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A Theory of Light Scattering in a Medlum With a Moving Boundary 78002 SOV/33-37-1-2/31

absorption coefficient per atom; n, the number of particles in a unit volume; x, a geometrical coordinate; $\mathcal{T} = \text{knx}$, the optical depth; t_1 , the average time a quantum is in a state of absorption; t_2 , the time spent by the quantum before two successive scatterings. Then \mathcal{T} may also be written as $\mathcal{T} = x/ct_2$ where c is the velocity of light. Two cases are considered: $t_1 \ge t_2$, and $t_2 \ge t_1$. In the first case, let $u = t/t_1$ be a dimensionless time, v, the velocity of the moving boundary, and $p(\mathcal{T}, u)$, the probability that a quantum of light absorbed at the depth \mathcal{T} will leave the medium in time t. Then if $P(\mathcal{T})$ is the probability of a quantum leaving the medium at any time, we have:

$$P(\tau) = \int_{0}^{\infty} p(\tau, u) du; \quad Z(\tau) = \int_{0}^{\infty} p(\tau, u) u du; \quad D(\tau) = \int_{0}^{\infty} p(\tau, u) u^{2} du. \quad (5)$$

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A Theory of Light Scattering in a Medium With a Moving Boundary $\begin{array}{l}
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\end{array}$ This integral equation is rewritten as: $P'(r) \coloneqq \frac{1}{2(1+r)} e^{-r} + \frac{\lambda}{2(1+r)} \int_{0}^{r} e^{-(r-r)} P(r') dr' + \\
 + \frac{\lambda}{2(1+r)} \int_{0}^{\infty} e^{-(r-r)} P(r') dr' - \frac{\lambda v}{1-r^{2}} \int_{0}^{\infty} e^{-\frac{r-r}{v}} P(r') dr', \quad (15)$ or $P(r) \coloneqq (1-k_{3}) e^{-kr}, \quad k_{3} \rightrightarrows \frac{1-\lambda}{v}. \quad (16)$ Here λ is an arbitrary constant. In the second case we $P(r) \equiv \frac{\lambda}{2} e^{-\frac{r}{1+v}} + \frac{\lambda}{2} \int_{0}^{\infty} e^{-(r-r)} p(r'-v|r-r'|) dr'. \quad (18)$ and $P(r) \equiv (1-k(1+r)) e^{-kr}, \quad k \equiv \frac{V4(1-\lambda)+\lambda^{3/2}}{2(1-r^{3})} = (2-\lambda)v. \quad (20)$ Card 3/4

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Some Notes on the Emission of Light under Cosmic Conditions

heating of the gas before the front of a shock wave can be calculated using the theory of light scattering in a medium with a moving boundary which was developed in Ref 5. In the one dimensional case, the intensity of radiation at an optical distance τ from the wave front is given by Eq (1), where $\tau = knx$, $v = Vknt_1$, k is the absorption coefficient per particle, n is the number of particles per cc, x is the distance from the wave front, V is the velocity of the wave front, λ is the ratio of the scattering coefficient to the total absorption coefficient (i.e. the sum of the true absorption and scattering coefficients) and t_1 is the mean lifetime of a quantum in the absorbed state. Eq (1) is subject to the conditions $(1 - \lambda) < 1, > 1$ which correspond to strong shock waves under cosmic conditions. The amount of radiant energy absorbed per unit volume and transformed into thermal

Card2/5 energy is given by Eq (2). As the volume element in the gas moves towards the shock wave-front, the energy

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Some Notes on the Emission of Light under Cosmic Conditions

accumulated in it is given by Eq (3), since dt/dx = -1/V. In a steady-state wave F, V and 1- λ retain unaltered. It then follows from Eq (3) that the energy $\frac{1}{2}$ is given by Eq (4), where $t_2 = 1/knc$ and is the mean lifetime of a quantum between two scattering events. In the first approximation one may put $F = \sigma T_{sh}$ in accordance with the Stefan-Boltzmann law where T_{sh} is the temperature on the front of the shock wave and is given by

 $T_{\rm sh} = \sqrt{3V^2/16R}$ where R is the gas

constant. For $1-\lambda$ the approximate relation is

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E032/E914

Some Notes on the Emission of Light under Cosmic Conditions

 $l-\lambda = \exp(-h\sqrt{kT_{sh}})$, where \sqrt{v} is the mean frequency of scattered radiation. A solution of the energy, mass and momentum conservation equations, which are given by Eq (5) with E given by Eq (4), determines the detailed structure of the heated region. It is, however, at once clear that the width of the heated region is approximately given by Eq (6). In <u>stellar</u> atmospheres this quantity is small and is of the order of a few centimeters or meters. In the chromosphere, the corona, or the interstellar gas, the width of the heated region is considerably greater and may become observable. Owing to the scattering of light in the higher-lying layers the radiation of the snock wave will penetrate into the outer layers before the shock wave reaches the surface. As a result, the intensity of radiation at the point of exit of the wave will begin to increase before the wave actually reaches this point. It is shown

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Akademiya nauk SSSR. Astronomicheskiy sovet		
Byulleton' stantsiy opticheskogo nablyudeniya iskusstvennykh sputnikov Zenli. no. 1 (11) (Academy of Sciences of the USSR. Astronomical Council. Bulletin of the Stations for Optical Observation of Artificial Earth Satellites. No. 1 (11)) Moscov, 1960. 22 p. 500 copies printed.		
Sponsoring Agency: Astronomicheskiy sovet Akademii nauk SSSR.		•
Resp. Ed.: Ye. Z. Gindin; Ed.: D. Ye. Shchegolov; Secretary: O.A. Severnaya.		-
PURPOSE: This bulletin is intended for scientists and engineers concerned with optical tracking of artificial satellites.		
COVERAGE: This bulletin contains short articles on optical equipment, techniques, and results of observations of artificial earth satellites. Also covered are the precision of satellite photography and the equations of motion of satellites. No personalities are mentioned. There are no references.		
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Observation Station	.I. Klimovskaya [L'vov]. On the Equation of Horizontal Coordinates			10		·	
Panaiotov, L.A. (M Artificial Earth Sa	ain Astronomical Observice tellites in the Polish	ratory]. Obse People's Repu	rvations of blic	12			- -
a) Bronkalla, b) Chuprina, A of the Astr	phic Observations of An V. Berlin-Babelaberg (.I., and L.A. Klepikovo consmical Council, AS U 1 Observatory	Deservatory a [Staff Hembe	h Satellites: TS	14 18			
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s/033/60/037/03/017/027 E032/E514

AUTHORS:	Kaplan, S.A. and Kurt, V.G.
TITLE:	On the Expansion of a Sodium Cloud in the Interstellar Space
PERIODICAL	: Astronomicheskiy zhurnal, 1960, Vol 37, Nr 3, pp 536-542 (USSR)
ABSTRACT: Card 1/4	Shklovskiy et al. (Refs 1 and 2) have described a method for the observation of the sodium cloud ejected from the second Soviet cosmic rocket on September 13, 1959. The results obtained by this method were also reported. The present paper gives a quantitative description of the expansion of the sodium cloud. It is shown that the expansion can be divided into two stages, namely, adiabatic expansion accompanied by a fall in the temperature and a free expansion during which the atoms preserve their thermal velocities corresponding to the temperature reached at the end of the adiabatic expansion. If one assumes spherical symmetry, then the expansion of the gas is described by Eq (3), where in the free expansion stage the term $\partial p/\partial r$ can be

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The second second of the second s S/033/60/037/03/017/027 E032/E514 On the Expansion of a Sodium Cloud in the Interstellar Space omitted. In the adiabatic stage the pressure gradient is also much smaller than the first two terms and the solution of Eq (3) is of the form given by Eq (4), where A is a constant and f(v) is an arbitrary function which is determined by the boundary and initial conditions. Certain hypothetical expressions for f(v) have been suggested by Stanyukovich (Ref 3). Under certain simplifying assumptions it can be shown that the relation between the velocity of adiabatic expansion and the thermal velocity of the second stage ck are related by Eq (7) in the case of spherical symmetry and by Eq (8) in the case of cylindrical symmetry. Assuming a Maxwell distribution of velocities (Eq 9), it is shown that the density distribution is given by Eq (12). Fig 1 shows the theoretical density distribution in the free expansion stage for various values of α which is proportional to the ratio a/c_k . Card 2/4 The dotted curve represents the density distribution

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\$/033/60/037/03/017/027 E032/E514 On the Expansion of a Sodium Cloud in the Interstellar Space form a new gas cloud which will expand with a lower velocity. The presence of such a secondary cloud may lead to a loss of definition of the central part of the main sodium cloud and to a slower fall off of the surface brightness. It is shown that this effect does not contribute appreciably to the outer structure of the main sodium cloud. Acknowledgment is made to L. M. Lukhovitskaya for assistance in the numerical computations. There are 3 figures, 1 table and 4 references, 3 of which are Soviet and 1 Dutch. ASSOCIATION: L'vovskaya astronomicheskaya observatoriya Gos. astronomicheskiy in-t imeni P. K. Shternberga (L'vov Astronomical Observatory, State Astronomical Institute imeni P. K. Shternberg) SUBMITTED: January 16, 1960 Card 4/4 /c **运动的环境的**部的

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s/033/60/037/005/005/024 E032/E514

AUTHORS :	Kaplan, 3.A. and Sivers, V.N.
TITLE:	The General Problem of Light Scattering in a One- Dimensional Medium with a Moving Boundary
PERIODICAL:	Astronomicheskiy zhurnal, 1960, Vol.37, No.5, pp. 824-827
$t_2 \rightarrow t_1$, wi absorbed sta between succ the general general equa quantum will of t_1 and one-dimensio probable in	In a previous paper (Ref.1), the authors investigated of the <u>scattering of light</u> in a one-dimensional medium mg boundary in the two special cases $t_1 \gg t_2$ and here t_1 is the lifetime of a light quantum in the ate and t_2 is the mean lifetime of the quantum cessive scatters. The present paper is concerned with solution of this problem and gives a solution of the ation for the probability that a scattered light leave the medium with a moving boundary for any values t_2 . As assumed before, the medium is taken to be onal and semi-infinite. The scattering is equally both directions and the probability of scattering is of the optical depth. The derivation is not given and
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	Problem of Light Scattering in a One-Dimensional Moving Boundary		
only the fina references.	l formulae obtained are quoted. There are 2 Soviet		×
ASSOCIATION:	L'vovskaya astronomicheskaya observatoriya (L'vov Astronomical Observatory)	/	Ð
SUBMITTED:	January 22, 1960		
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Effect of anisotropic conductivity in a magnetic field on the structure of a shook in magnetic gas dynamics. Zhur. eksp. 1 teor. fiz. 38 no.1:252-253 Jan '60. (MIRA 14: (MIRA 14:9)

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1. L'vovskiy gosudarstvennyy universitet. (Magnetic fields) (Shock waves)





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3/702/62/000/009/001/002 1046/1246 AUTHOR: Kaplan, J.A. TITLE: The determination of the optimal excitation conditions of elastic SOURCE: USBR. Glavnoye upravleniye geologii i okhrany nedr. Geofizicheskaya razvedka, no. 9, 1962, 28-36 . . The conditions of excitation are assessed from the amplitudes of the reflected TEXT: waves generated in microseismotorpedoing. This method cannot be used unless a) discontinuities with high reflection coefficients exist within the seismogeological cross V section, and b) the reflected base waves are known to require identical or similar excitation conditions. There are 3 figures. Card 1/1

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AUTHORS:	Kaplan, S. A., Kutik, I. N.		
TITLE:	On the emission of magnetohydrodynamic and magnetoaccust c waves		1 1 1 1
PERIODICAL	Referativnyy shurnal, Fisika, no. 3, 1963, 6, abstract 30-1 ("Visnyk L'vivak'k. un-tu. Ser. fiz.", 1962, no. 1(8), 75 - 78, Ukrainian)		
eren en en en en en		计过去分词分子	
TECT:	The propagation of magnetohydrodynamic waves is considered in the		
case where Solving the	The propagation of magnetohydrodynamic waves is considered in the the source of oscillations is expressed in the form $F = 22^{-12} e^{-110/4}$, equations of magnetic hydrodynamics, the authors obtain an expression		
case where Solving the for magnete	The propagation of magnetohydrodynamic waves is considered in the the source of oscillations is expressed in the form $F = 22^{-27}/a^{-110/4}$.		
case where Solving the for magnete	The propagation of magnetohydrodynamic waves is considered in the the source of oscillations is expressed in the form $F = 2^{2\pi/3} + 10^{4}$, equations of magnetic hydrodynamics, the authors obtain an expression bydrodynamic and magnetoscoustic waves. Expressions are obtained for		
case where Solving the for magnete the average	The propagation of magnetohydrodynamic waves is considered in the the source of oscillations is expressed in the form $F = 2^{2\pi/3}/4^{+10/4}$, equations of magnetic hydrodynamics, the authors obtain an expression hydrodynamic and magnetoscoustic waves. Expressions are obtained for doing intensities of the emission of the mentioned waves.		
case where Solving the for magnete the average	The propagation of magnetohydrodynamic waves is considered in the the source of oscillations is expressed in the form $F = 2^{-1/2} \cdot 2^{-1/2} \cdot 10^{-1}$, equations of magnetic hydrodynamics, the authors obtain an expression bydrodynamic and magnetomoustic waves. Expressions are obtained for d-in-time intensities of the emission of the mentioned waves.		
case where Solving the for magnete the average	The propagation of magnetohydrodynamic waves is considered in the the source of oscillations is expressed in the form $F = 2^{-1/2} \cdot 2^{-1/2} \cdot 10^{-1}$, equations of magnetic hydrodynamics, the authors obtain an expression bydrodynamic and magnetomoustic waves. Expressions are obtained for d-in-time intensities of the emission of the mentioned waves.		

CIA-RDP86-00513R000520430013-3

\$/058/63/000/c03/028/104 A062/A101 AUTHORS : Kaplan, S. A., Koval Chuk, V. G., Korolishin, V. M. Coefficients of electric conductivity and diffusion in relativistic TITLE: one-component plasma Referativnyy Mulmal, Fisika, no. 3, 1963, 19, abstract 3113 PERIODICAL ("Vienyk L'vive"k, un bu, Ber, fis.", 196?, no. 1(8), 79 + 82, Ukrainian) A method is given for computing the coefficients of diffusion and TEXT: electric conductivity in a relativistic one-component plasma in the presence of electric and magnetic fields. Expressions for the components of the "fourdimensional velocity" of the particles are averaged, for the cases of parallel and perpendicular electric and magnetic fields, by means of the distribution function in the zero approximation. Transfer coefficient is obtained in the presence of an electric field and the gradient of concentration of the particle. For a relativistic plasma, at a power exponent of the particle spectrum $\gamma = 2$, the diffusion coefficient is inversely proportional to the intensity of the magnetic field. [Abstracter's note: Complete translation] Yu. Mordvinov Card 1/1

APPROVED FOR RELEASE: 06/13/2000

ACCESSION NR: APLO07673

s/0214/63/000/006/0053/0059

AUTHORS: Kaplan, S. A.; Ostrovskiy, L. A.

TITLE: Theory of shock wave formation in chromosphere and corona

SOURCE: Solnechnywye dannywye, no. 6, 1963, 53-59

TOPIC TAGS: acoustical theory, sound wave, sound velocity, magnetic force tube, energy dissipation, shock wave, coronal shock wave, supersonic flow, gas flow, corona, chromosphere, wave formation

ABSTRACT: The authors have examined the conditions for converting sound waves to shock waves in an inhomogeneous atmosphere within a gravitational field. This consideration is associated with determination of magnetic turbulence. The authors describe the application of a method that permits investigation of conditions for converting sound waves to shock waves in any distribution of density and temperature, under conditions that the wave length of the sound is much less than the equivalent height and that self-excitation is small. The method has been discussed elsewhere by K. Ye. Gubkin (Sb. "Nekotorywye problemyw matematiki i mekhaniki" AN SSSR, Novosibirsk, 1961, str. 69) and 0. S. Ry*shov (Zh. prikl. mekh. i tekh. fiz.,

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APPROVED FOR RELEASE: 06/13/2000

CIA-RDP86-00513R000520430013-3"

ACCESSION NR: AP4007673

no. 2, 15, 1961). The authors consider velocity of the gas, the effect of gravity, and energy flux. From the relationship that shock waves form when the steepness of the sound-wave front approaches infinity, they find expressions for the distance a sound wave must travel before rupture occurs (that is, before a shock wave is generated). This distance is found to be on the order of 10° cm. The distance a sound wave will travel before half its energy is dissipated is on the order of $2 \cdot 10^{\circ}$ cm. It is concluded that a substantial part of the kinetic energy of the chromosphere. It is possible that this circumstance explains the sharp rise in temperature at the inner boundary of the corona. Further dissipation of energy a high temperature gradient. Orig. art. has: 30 formulas.

ASSOCIATION: Gor'kovskiy nauchno-issledovatel'skiy radiofizicheskiy institut (Gorkiy Scientific Research Radio Physics Institute)

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AUTHOI	1: Kaplan, B.A.			<u>يونيد ميرونيوني</u> ت	·6	
TITLE:	Some problems of	the physics of the inter	stellar and inte	rplanetary mode	i i - i - 1 - 1 - 1 - 1 - 1 - 1 - 1 - 1 - 1 -	:
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ABSTRA(CT: This three-pa is calculated using B. Pikol'ner, Me and Milligan (An	rt paper deals with 1) a method described in ezhzvezdnaya sreda, M J., 137 1, 1962) K wayo metari	Interstellar ra	r publication is d	1	
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the motion of such a wave (and the subsequent pressure increases) are calculated using the method of self-similar motion of ionization explosions (outlined in §15 of S.A. Kaplan, Mezhzvednaya gazodinamika, M., Fizmatgiz, 19:8); and 3) the transfer of radiation within a spherically symmetrical darkened space (he problem is solved in the Eddington approximation assuming that the density of the scattering medium decreases (from the center) as 1/r³. In a subsequent paper, these calculations will be applied to the scattering is of L₄ in the night glow spectrum. Orig. art. hus 25 formulas. 2 figures. FN : correctly instruments as formula and the state of the scattering medium decreases (from the center) as 1/r³. In a subsequent paper, these calculations will be applied to the scattering is of L₄ in the night glow spectrum. Orig. art. hus 25 formulas. 2 figures. FN : correctly instruments as

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"APPROVED FOR RELEASE: 06/13/2000 CIA-RDP86-00513R000520430013-3 KAPLAN, S.A.; ZAYTSEV, V.V.; KISLYAKOV, A.G.; KOBRIN, M.M.; TSEYTLIN, N.M. Fourth All-Union Conference on Radio Astronomy. Izv. vys. ucheb. zav.; radiofiz. 6 no.4:861-869 '63. (MIRA 16: (MIRA 16:12) 1.1.1.1 The set and a sub-

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KAPLAN, S.A.

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Spectrum of magnetohydrodynamic turbulent convection. Astron. zhur. 40 no.6:1047-1054 N-D '63. (MIRA 16:12)

1. Radiofizicheskiy institut Gor'kovskogo gosudarstvennogo univorsiteta.

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"Measurements of scattered U. V. radiation (1216A and 1300A) in the upper atmosphere"(USSR)

Report submitted for the COSPAR Fifth International Space Science Symposium, Florence, Italy, 8-20 May 1964.

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AP4015565	t	S/0089/64/016/002/0149/0150	
AUTHOR: Zaytsev	, V. V.; Kaplan, S. A	A.	
TITLE: Concerni scattering of g	ng the theory of the amma photons	nonstationary multiple Compton	
SOURCE: Atomnay	a energiya, v. 16, no	o. 2, 1964, 149-150	
TOPIC TAGS: mul photon, Compton	tiple Compton scatter scattering	ring, small angle, photon, gamma	
ABSTRACT: This nonstationary s approximation	paper presents a simp cattering of gamma ph	ele solution of the problem of the otons for small angles. The	28
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KAPLAN, S.A.; KLIMISHIN, I.A.

Methods of analysis of interstellar turbulence. Astron.zhur. 41 no.2:274-281 Mr-Ap '64. (MIRA 17:4)

1. L'vovskaya astronomicheskaya observatoriya i Radiofizicheskiy institut Gor'kovskogo gosudarstvennogo universiteta.

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CIA-RDP86-00513R000520430013-3

ACCESSION NR: AP4043953

S/0033/64/041/004/0652/0656

AUTHOR: Dibay, E. A., Kaplan, S. A.

TITLE: Cumulative shock waves in interstellar space

SOURCE: Astronomicheskiy zhurnal, v. 41, no. 4, 1964, 652-656

TOPIC TAGS: astrophysics, interstellar space, shock wave, cumulative shock wave, interstellar gas, globule, star, nebula, Stromgren zone

ABSTRACT: Dense circular dust nebulae (globules) are frequently observed within H II emission regions. As a result of the sharp temperature difference between the globule , and the surrounding ionized medium it is possible to expect its compression by a shock wave developing at the discontinuity. If the configuration of the globule is close to spherical the shock wave will have a cumulative character, that is, there will be focussing of the wave toward the center. If a dark nebula in a H II zone is greatly elongated it is also possible to have cylindrical cumulation. At the time of development of a type 0 star, causing the ionization of a surrounding nebula, a Strömgren zone is formed around it. If there are such dense fluctuations within the nebula that it cannot be penetrated by ionizing radiation, the H II zone will "bend around" such formations. The time required for establishment of the Strömgren zone is of the order of the time required for recombination Cord 1/3

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ASSOCIATION: Gosuda (State Astronomical Ins universiteta (Radiophys	incute): Kad	llofizichesk	tiv institu	nstitut imeni P. K. Shternberga t Gor'kovskogo gosudarstvennogo tiversity)	i
has: 25 formulas.	s can de obi	tained for a	. cylindric	where between the values cited al cumulative wave. Orig. Art.	•
Shock wave velocity	U _d Nd	220 1.1	10 ⁵ 2.4	Reflected shock wave Ditto	
Gas velocity Density		1.05	2.36	Ditto	
Shock wave velocity Density	71 _ธ บ	0.45 10	0.75 10 ³	Ditto Attains center	:
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un de la companya de L 7048-65 0 ACCESSION NR1 AP4043954 itial velocity of the emission front, $n = o_0/\rho$, where o_0 is density of the flowing gas and p is the density of the nondisturbed gas, y is the feentropic exponent, R is an ideal gas constant, and u is the nolacvlar weight. On the basis of light scattering theory, differential equations for the source function B(1) are written and helt solutions containing terms T+ (temperature behind the shock wave front) and T_ (temperature ahead of the shock wave front) are derived. To detorminat the unknowns T_{+} and T_{-} , expressions derived earlier for T and F for the values behind and shead of the shock wave front are used. a set theor equations is derived from which is an o are calcul to a n even ators and press. - established that th the distance from the 4 The results obtained are so . 4 ¥ . propagation of shock waves in stellar envelopes. IT IS SNOVE THAT remeasures in the heated some immediately ahead of and behind the shock wave front are of the same order. Orig. art. has: 14 formulas. ard 113



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AUTHOR: Kaplan, S. A		2
TITLE: Relativistic	shock waves in intergalactic ges	
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ACCES	TON NR: AP2009640	.; Karpinskiy, I. P. ; Pustoveyt, R. M.;	<u>UN/V293/0</u>	01003100610631		
TITLE	: Investigation of	scattered ultraviole				
	quipment 5: Kosmicheskiye is	sledovaniya, v. 3, n	10. 2, 1965, 237-243	3	r	
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ACCESSION NR: AP5009643 UR0293/65/003/002/0251/0256	1.7
AUTHOR: Kaplan, S. A.; Kurk, V. G.	
TITIE: The theory of the resonance scattering of L sub AL, is redistion in the	
SOURCE: Kosmicheskiye issledovaniya; v. 3, no. 2, 1965, 251-256	
TOPIC TACS: geocorona, L sub Alpha radiation, resonance radiation theory, upper atmosphere, hydrogen distribution	
IPSTRACT: The authors point out that the various papers which have appeared in re 	r
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KAPLAN, S.A.

Continuous salt method for the production of synthetic nitron fibers. Biul. tekh.-ekon. inform. Gos. nauch.-issl. inst. nauch. i tekh. inform. 18 no.3:20-21 Mr '65. (MIRA 18:5)

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KAPLAN, S.A.; PETRUKHIN, N.S.

Theory of convection in a polytropic atmosphere with a uniform magnetic field. Astron. zhur. 42 no.1:74-77 Ja-F '65.

(MIRA 18:2)

1. Radiofizicheskiy institut Gcr'kovskogo gosudarstvennogo universiteta i Ural'skiy gosudarstvennyy universitet.

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KAPLAN, U.S.; KURT, V.G.

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Theory of the resonance scattering of 1 -radiation in the geocorona, Kosm.1ssl. 3 no.2:251-256 Mr- 165.

Interpretation of observations of the triplet OI (1300 Å) in the upper atmosphere. Ibid.:256-261 (MIRA 18:4)

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KAPLAN, S.A.; LUPANOV, G.A.

Relativistic instability of polytropic spheres. Astron.zhur. 42 no.2:299-304 Mr-Ap *65. (MIRA 18:4)

1. Gor*kovskiy nauchno-issledovatel*skiy radiofizicheskiy institut.

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ACCESSION NR: AT502357	72	en e		
AUTHORS: Kaplan, S. A.	s Kurt, V. G.		/000/000/0111/0112	801
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TOPIC TAGS: solar radia phere, atmosphere model,	tion scattering, sol	ar radiation absorpt	tion, upper stars	
ABSTRACT: The scatterin of the earth is consider great optical thickness. to take place in the upp oxygen occurs in the low	g of ΟΙ(λ1300 Å ed, using the double Scattoring of the j er layer without abso or layer where it is) radiation in the u layer model of the incident solar radia orption. Absorption	pper atmosphere atmosphere of tion is assumed by molecular	
scattering event of \wedge in solution of the shift equ that the intensity begins which agrees well with ob Cord $1/2$	ation spplied to the	h increasing optical	L thickness The	
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	SOURCE CODE: UR/0269/65/000/010/0042/0042
AUTHORS: Kaplan, S. A.; Petrukh	<u>in, N. S.</u> 29
TITLE: Interpretation of the "s photosphere	upersonic" propagation of disturbances in the solar
SOURCE: Ref. zh. Astronomiya, A	be. 10.51.311
REF SOURCE: Solnechnyye dannyye	, no. 10, 1964(1965), 63-66
	solar disturbance, solar magnetic field
with a velocity up to 2 km/sec wa an intensity up to 100 oe located ing some decrease of the magnetic which the gas began to ascend at central part of the magnetic hill along the surface of the sun to t westup to 280 km/sec. The auth ward drift of a magnetic force to	etation is given of the phenomenon observed by G. aph of GAO 20 July 1961. A sharp descent of gas as observed in the region of a magnetic hill with d far from sunspots. This descent occurred follow- c hill intensity and lasted about 1/, min, after half the velocity. The descent of gas began in the l, then the front of the region began to propagate the east with a velocity of 50 km/sec and to the hors assume that the phenomenon began with the down- ube originally located at a fixed depth z ₀ , because
of which a zone of variable dist	urbance originated in this region. Sonic dilatation
ard 1/2	UDC: 523.74
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"supersonic" propagation of the g calculation of the propagation to sun is carried out; the distance to the point of ray emergence is between the ray direction and the that the solar atmosphere is poly with depth. It is shown that for of the emergence point of sonic ways at the depth of the so sonic ways at the depth of the so scent zone front of 50 km/sec, the tude of the disturbance propagation of the assumption that the sonic the depth z_0 , but extends at this to the horizontal. Thereby it is	SE: 06/13/2000 CIA-RDP86-00513R00052043D sides from this zone. The subsequent emergence to at various angles to the normal led to the observed gas descent zone along the surface of the sum. A ime of sonic dilatation waves to the surface of the along the surface from the point above the source also calculated as a function of the angle φ_0 e surface normal. It is assumed for the calculation ytropic and that the temperature gradient is constant $f \varphi_0$, not too close to zero, the velocity of motion saves to the surface is close to the velocity of source. For a propagation velocity of the gas de- te source depth $s_0 = 20 000$ km. The greater magni- ton velocity to the west is explained on the basis wave source is not concentrated in a small volume at a depth in the latitudinal direction at a small angle assumed that the region of the original magnetic



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