ZHELUDEY, I. S., ABDULGAMIDOV, I. S., FRIDKIN, V. M.,

"Testing of photoelectret conditions in activating alkali halogen crystals."

report to be submitted for the lat Intl. Congress on Reprography, Cologne,

West Germany, 14-19 Oct 1963

ZHELUDEV, I. S., YURIN, V. A.,

"Stabilization of Spontaneous Polarization in Ferroelectric Crystals."

report presented at the Symposium on Ferroelectricity and Ferromagnetism, Leningrad, 30 May-5 June 1963.

45676

24,7000

\$/070/63/008/001/008/0841 E132/E460 - 7

AUTHORS:

Vlokh, O.G., Zheludev, I.S., Shamburov. V.A.

TITLE:

The electro-optical effect in crystals of penta-

erythritol C(CH2OH) 4 ..., nut 1.

PERIODICAL: Kristallografiya, v.8, no.1, 1963, 51-56

TEXT: For pentaerythritol, which belongs to the crystal class 4, crystals showing the growth pyramids 100 appear to have a two-fold axis. Abstracter's note: The authors state that the crystals appear biaxial optically. This does not appear to be correct as this system must be uniaxial, but it may mean that the ellipsoid of revolution which represents the refractive indices requires two parameters to describe it and has two different axes. The optical indicatrix is described by the equation:

$$(a_0^2 + r_{12}E_y)x^2 + (b_0^2 + r_{22}E_y)y^2 + (c_0^2 + r_{32}E_y)z^2 + 2r_{52}E_yzx = 1$$

when an electric field Ey is applied along the y-axis. This y-axis is the fourfold inversion axis for the crystal as a whole. ao, bo and co are the reciprocals of the principal Card 1/3

s/070/63/008/001/008/024 E132/E460 The electro-optical are the electro-optical coefficients. refractive indices; r. 8 being non-zero for this cut. It follows that when an electric field is applied the indicatrix is deformed and rotates in the XZ plane through an angle 52. This y-cut crystal was mounted between crossed Nicols and a beam of monochromatic light was passed through the system into a photomultiplier. The plate was adjusted to extinction and a high voltage was applied to the electrodes, the increase in transmitted light being measured. The increase resulted from the rotation of the indicatrix which could reach 22.5° if a field of 220 kV/cm were applied. material has a high melting point (257°C) and behaves as a linear dielectric with a specific resistance of  $10^{15}$  to  $10^{12}$  ohm cm over the range 30 to 130°C in the absence of surface conductivity. The crystals are not hygroscopic and have a perfect 001 cleavage which corresponds to the y-cut used if it is reckoned that the growth pyramids of the form 101- give crystals of the class 2. The moduli were found to be  $r_{52} = (4.38 \pm 0.13) \times 10^{-8} \text{cgsu}$ Card 2/3

S/070/63/008/001/008/024

The electro-optical ...

The rotation of the indicatrix is about 20 times that produced by the same field in ammonium dihydrogen phosphate. The latter (ADP) has, however, a much greater electro-optical effect when the field is in the Z-direction. There are 3 figures.

ASSOCIATION: Institut kristallografii AN SSSR (Institute of Crystallography AS USSR)

SUBMITTED: June 26, 1962

45677

24,2130 24,7800

8/070/63/008/001/009/024 E132/E460

AUTHORS:

Sonin, A.S., Zheludev, I.S.

TITLE:

The investigation of certain dielectric properties of

single crystals of sodium nitrite

PERIODICAL: Kristallografiya, v.8, no.1, 1963, 57-62

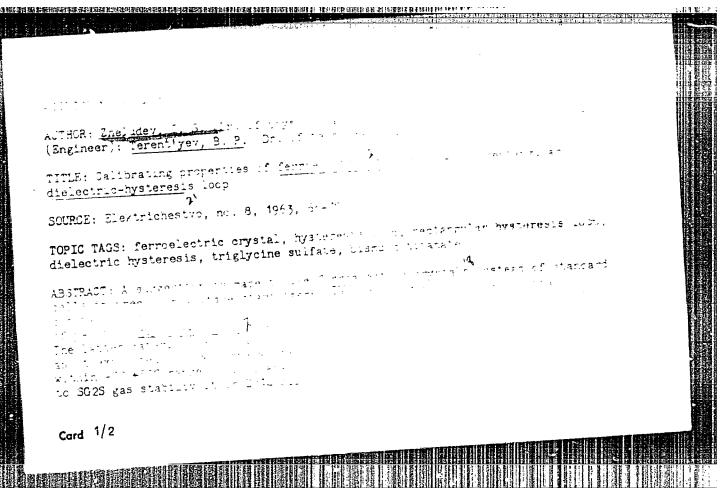
TEXT: Ferroelectricity in NaNO<sub>2</sub> was predicted by the authors in 1957. Crystals grown from solution were used first but were unsatisfactory because of high electrical conductivity. The dependence on temperature of dielectric constants, spontaneous polarization, coercive force, electric conductivity and piese-electric properties of single crystals of sodium nitrite were investigated. The crystals used were made from the melt in closed ampules from salt which had been earlier dried in vacuo at a high temperature. They were non-hygroscopic. They, cleaved easily parallel to (101). Plates of other orientations were thus difficult to prepare. Electrodes were applied by evaporating silver. The dielectric constants were measured at 500 Kc/s and, at room temperature, were e<sub>a</sub> = 6.8, e<sub>b</sub> = 6.4 and e<sub>c</sub> = 7.8; at 170°C a λ-point was found in the curve for e<sub>c</sub> Card 1/2

### S/070/63/008/001/009/024 E132/E460

The investigation of certain ...

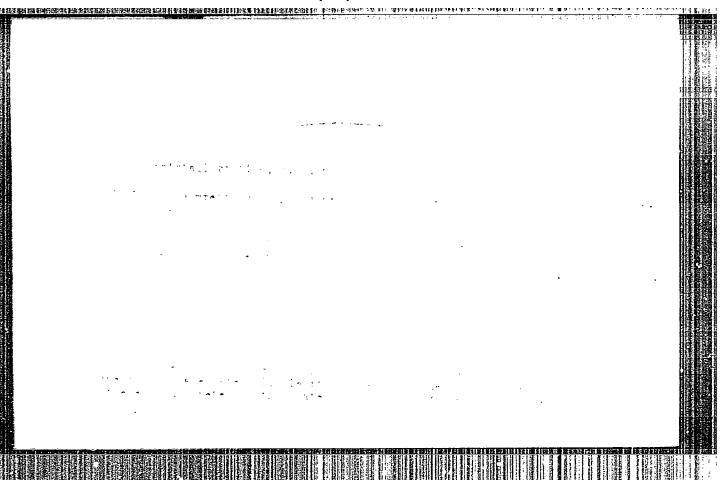
which rose to above 800. Above the Curie point of 161°C ec obeys the Curie-Weiss law. ea and eb show discontinuities of slope at the Curie point but do not exceed 14. The hystoresis loop was plotted at various temperatures. At room temperature the coercive force is higher than could be applied. Above 120°C the hysteresis loop could be observed. The spontaneous polarization was measured as 8 microcoulombs/cm² and at 150°C the coercive force was 2 kV/cm changing linearly with temperature. The piezo-electric moduli were measured, d33 being 4.4 x 10-8 cgsu at room temperature. It changed little until above 160°C when a rapid fall to zero occurred. The electrical conductivity followed the law, log s = k/T, giving an energy of activation of 0.72 eV above the Curie point and 0.90 in the ferroelectric region. Nonlinear effects were found when the susceptibility was measured as a function of field at 1 Kc/s. The behaviour of NaNO2 is to be compared with that of KNO3 where there is also no hydrogen. The ferroelectricity and the transition appears to be connected with ordering of the NO2 groups. There are 9 figures.

ASSOCIATION: Institut kristallografii AN SSSR (Institute of SUBMITED: Pay 5, 1962 Crystallography AS USSR). Card 2/2

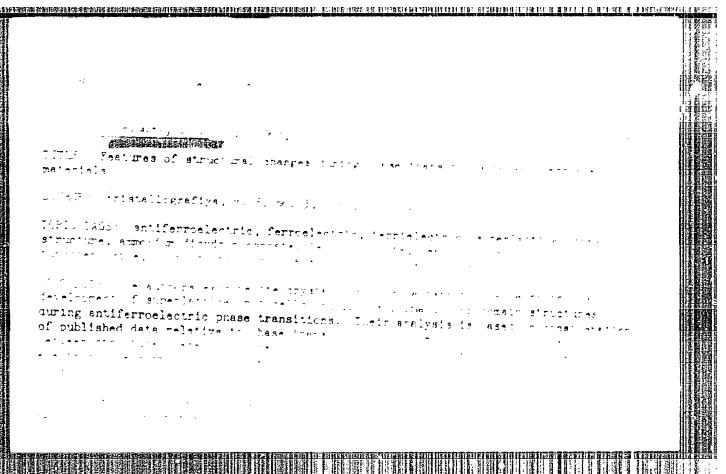


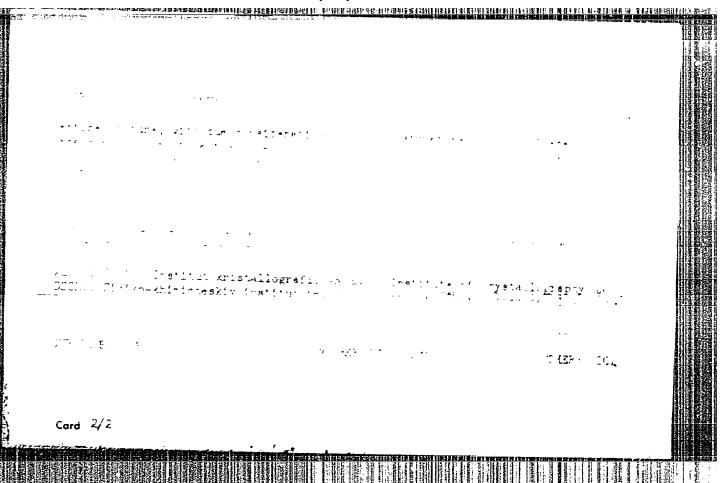
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TAMBOV	SEV, D.A.; SKORIKOV, V.M.; ZHELUDEV, I.S.	
	Production of bismuth titanate single crystals and their properties. Kristallografiia 8 no.6:889-89	93 N-D'63.
	1. Institut kristallografii AN SSSR.	(MIRA 17:2)
	마음보다 그렇게 되었다. 그 생각이 되는 사람이 되는 것이 되었다. 그 사람들이 되었다. - 1일 부분에 가지 하는 것이 보면 하는 것이 되었다. 그 사람들이 되었다. 그 사람들이 되었다.	
	환기 가장하다는 한국장에 불안하면서 그렇게 되었다. 그 사람이 있는데 한 역사 (2010년) 전화 경영화 기관을 다 있는 것이 되었다. 기계	
	그는 말았다면서 문화하면 중에 등을 통해 보고 있습니다. 그런 그는 것이 되어 생각을 받는다. 기본 경기 교통 전기 교통 전기 기본	
	. 물리 후에 내 다른 명을 보고 한 생활되는 이 생활 경기에는 살았다고 되었다. 1945년 - 이 대한 회원 회원회사 대한 학생 등 보고 기업을 받았다.	일 제한 12 명이 되고 해를 받는 역사 전략으로 가는 스트를 받는다.
	게 즐리고 있는 것 같아. 그런 그런 그런 그런 그런 그런 그런 것이다. 사용에 살아왔다. 그런 그런 그런 그런 그런 그런 그런 것 같아.	
로 보이스러 누를 받는다. 보이트 보기 하는데	- 발생하게 보는 사람이 되었다면 하는 것이 되었다. 그 사람이 되었다. - 발표를 하는 사람들이 들었다면 있다. 그는 사람들이 되었다.	
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	사용하는 경기를 보고 있다. 그는 말을 보고 있는 것이 되었다. 사용물 및 물을 보고 말을 보는 것을 보고 있습니다. 그 것이 되었다. 그는 것이 없는 것이다.	

## S/0105/64/000/003/0001/0005

ACCESSION NR: AP4019324

AUTHOR: Tambovtsev, D. A. (Engineer); Terent'yev, B. P. (Doctor of technical sciences); Zheludev. I. S. (Doctor of physico-mathematical sciences);

Skorikov, V. M. (Engineer); Kucherova, I. V. (Engineer)

TITLE: Voltage and current stabilization by ferroelectrics

SOURCE: Elektrichestvo, no. 3, 1964, 1-5

TOPIC TAGS: ferroelectric, ferroelectric crystal, voltage stabilizer, current stabilizer, ferroelectric voltage stabilizer, ferroelectric current stabilizer, reference voltage, bismuth titanate, barium titanate, triglycine sulfate

ABSTRACT: Procedures for the calculation of ferroelectric-stabilized reference-voltage sources are set forth, a new circuit for voltage stabilization is submitted, and some problems in using ferroelectrics for stabilization purposes are discussed. The new bridge-like circuit (see Enclosure 1) has the advantage

Card 1/32

CIA-RDP86-00513R002064710012-8" APPROVED FOR RELEASE: 03/15/2001

ACCESSION NR: AP4019324

of a high output voltage that can reach one-third of the input voltage; also, a high degree of temperature compensation is possible. The experimentally determined effects of frequency and load on the performance of ferroelectric voltage stabilizers are reported. The possibilities of ferroelectric materials for current stabilization were also explored; a 1-cm<sup>2</sup> barium-titanate plate ensured a stable mean current of 50 ma at 50 cps; bismuth titanate and triglycine sulfate were also tested. Orig. art. has: 9 figures and 6 formulas.

ASSOCIATION: Institut kristallografii AN SSSR (Institute of Crystallography,

SUBMITTED: 13Sep63

DATE ACQ: 27Mar64

ENCL: 01

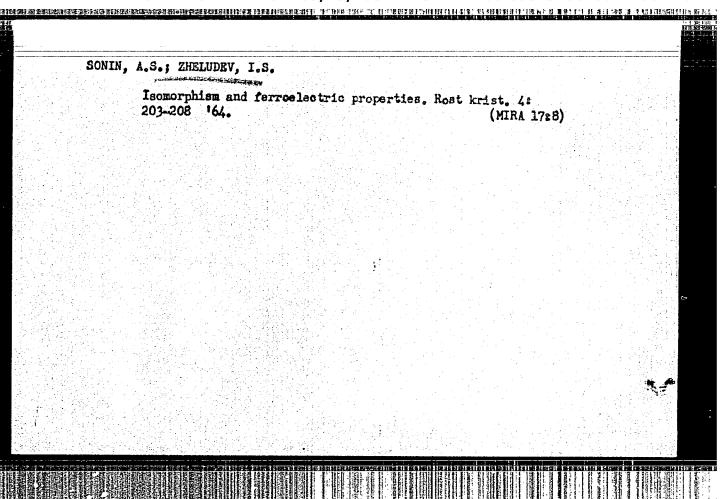
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APPROVED FOR RELEASE: 03/15/2001 CIA-RDP86-00513R002064710012-8"



ACCESSION NR: AP4019835

8/0181/64/006/003/0764/0770

AVIHORS: Abdulgamidov, S. A.; Zheludev, I. S.; Nosov, V. N.; Fridkin, V. M.

TITIE: On internal field distribution in single crystal photoelectrets

SOURCE: Fizika tverdogo tela, v. 6, no. 3, 1964, 764-770

TOPIC TAGS: internal field distribution, single crystal, photoelectrot, photoelectric field, interelectrode spacing, space charge, field distribution

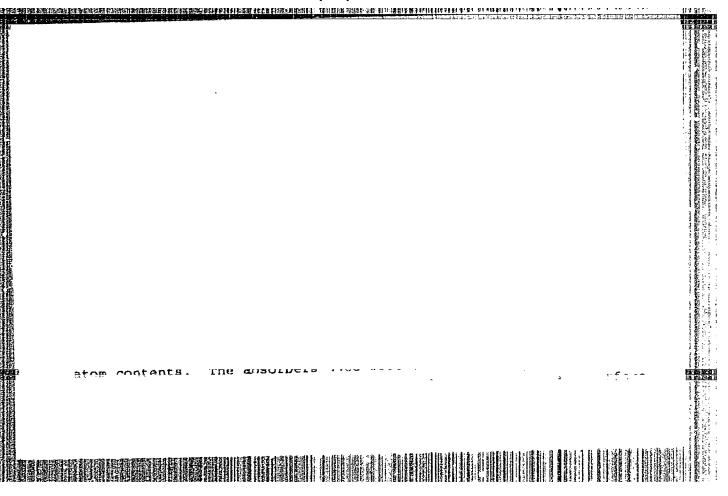
ABSTRACT: The hetero- and homocharge distributions in photoelectrets of single crystals of additive-colored KCl, S, K2Cr2O7 and CdS have been investigated, using the light probe technique of M. Y. Ben Sira, B. Pratt, E. Harnik, and A. Many (Phys. Rev. 115, 55, 1929) and Harnik, Ben Sira, Pratt, and S. Peter (J. Appl. Phys., 34,207, 1963). It consists of depolarizing the photoelectret by means of a light probe in a direction perpendicular to the internal photoelectric field. The KCl specimen was polarized first by a 0.5-kv field with exposure of the whole crystal. surface to 546m monochromatic light, and subsequently by a 2.0-kv field with central illumination only. Internal field distributions are represented graphically as functions of the interelectrode spacing. Both barrier type and space

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and normal property in the control of the control o

## ACCESSION NR: AP4019835 charge type distributions are observed. For the central illumination case an asymmetric field distribution was noticed relative to the crystal center. Similar experiments were performed on the rest of the specimens. In CdS, under all polarization time durations, the field showed an inverse direction at the cathode and a forward, positive direction at the anode. Orig. art. has: 6 figures and 3 formulas. ASSOCIATION: Institut kristallografii AN SSSR Moscow (Institute of Crystallography AN SSSR) SUBMITTED: 29Aug63 DATE ACQ: 31Mar64 ENCL: 00 SUB CODE: PH NO REF SOV: 006 OTHER: 006

Card 2/2





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	1. Institut kristallografii AN SSS	ir.	(MIRA 17:11)
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s/0070/64/009/004/0564/0565

ACCESSION NR: AP4043192

AUTHORS: Abdulgamidov, S. A.; Zheludev, I. S.; Fridkin, V. M.

TITLE: On the kinetics of the formation of the photelectret state in additively colored KCl single crystals

SOURCE: Kristallografiya, v. 9, no. 4, 1964, 564-565

TOPIC TAGS: potassium compound, photoelectret, single crystal, color center, F center, polarization, depolarization

ABSTRACT: The tests were made with single crystals having an F-center concentration  $\sim 10^{16}$  cm<sup>-3</sup> with an aim at checking on the theoretical deductions of E. Adirovich (Fiz. tv. tela v. 3, no. 7, 2048, 2050, 1961). The measurements were made with plates having area  $\sim 1-2$  cm<sup>2</sup> and thickness  $\sim 0.15$  cm, using barrier contacts to prevent the injection of the carriers from the electrodes into the

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ACCESSION NR: AP4043192

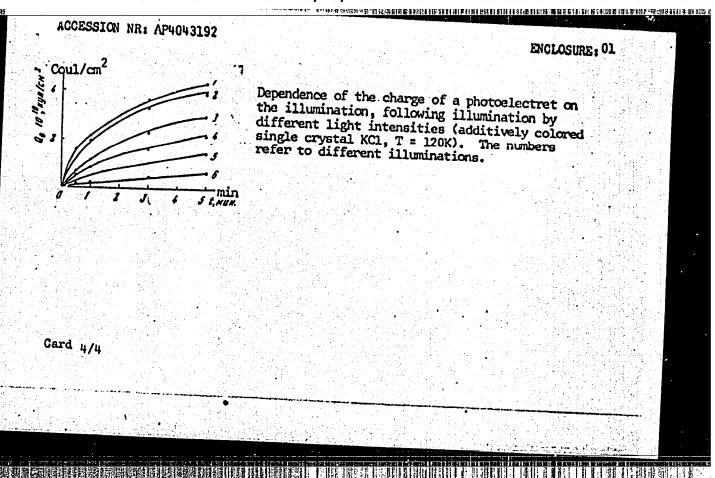
crystal (mica liners). Illumination during the polarization was with monochromatic 546-nm light. The charge was measured by the depolarization method, with the photoelectret illuminated with the undecomposed light from a mercury lamp and with integration of the depolarization current. The time variation of the photoelectret charge was found to be exponential. The reciprocity law was fulfilled over the entire range of light intensity employed. The charge density was found to be linear in the polarizing-field intensity. It is concluded on the basis of these results, and also results by others, that the kinetics of formation of the photoelectret state in colored alkali-halide crystals, as well as their depolarization, are in accord with the deductions of the phenomenological theory. "The authors thank L. M. Shamovskiy for supplying the samples of additively colored single crystals." Orig. art. has: 1 figure.

ASSOCIATION: Institut kristallografii AN SSSR (Institute of

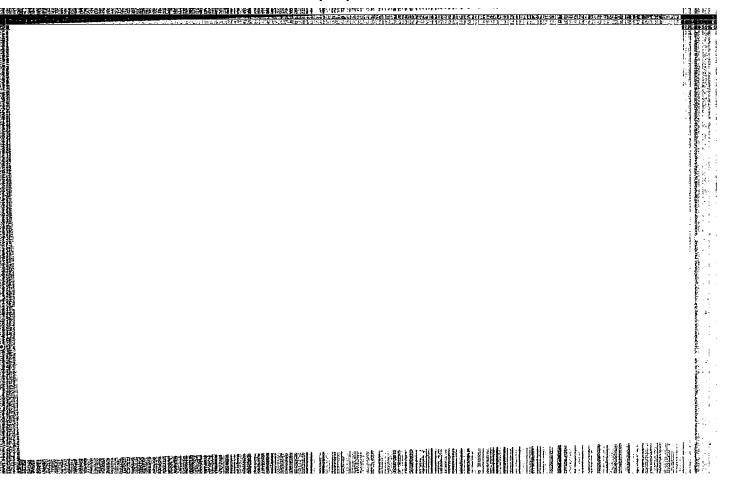
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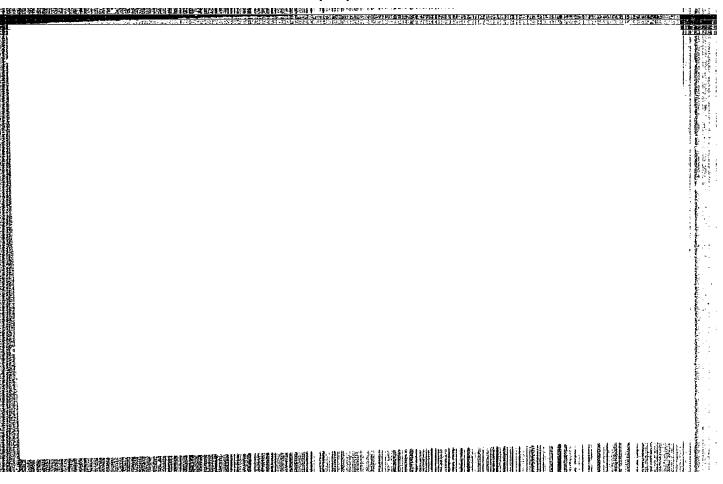
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Crystallography, AN SSSR)
SUBMITTED: 01Nov63 ENCL: 01
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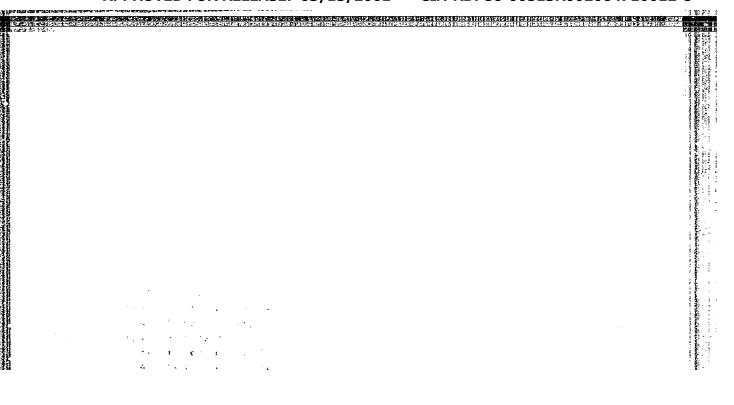
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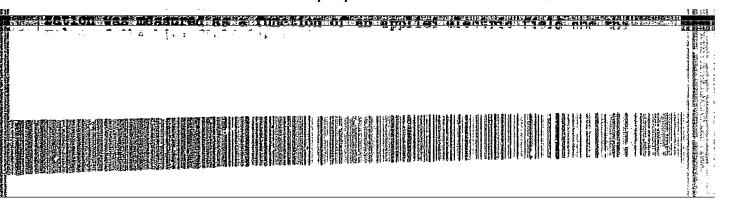


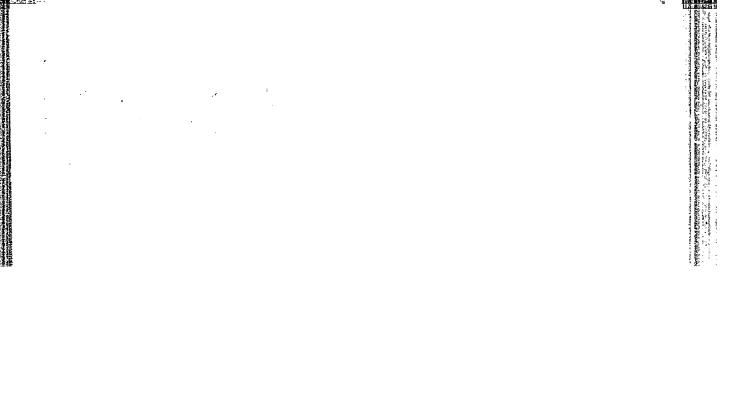
# ROMANYUK, N.A.; ZHELUDEV, I.S. Changes in the dome'n structure of Rochelle salt crystals due to radiation. Kristallografija 9 no.6:876-878 N-D '64. (MIRA 18:2) 1. L'vovskiy gosudarstvennyy universitet i Institut 'ristallografii AN SSSR.

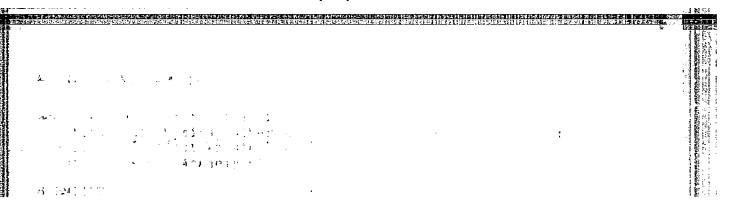


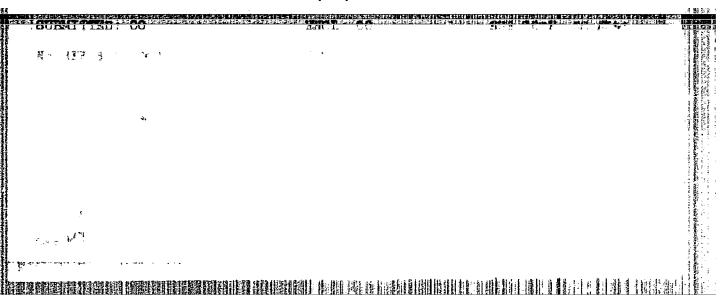












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L 4271-66 ENT(1)/EPA(s)-2/ENT(m)/ENP(t)/ENP(b) IJP(c) 1.1 JD ACCESSION NR: AP5024569 UR/0070/65/010/005/0764/0766 548.0:537.226.1

AUTHOR: Zheludev, I. S.; Ludupov, Ts-Zh.

TITLE: Complex dielectric constant of RbH<sub>2</sub>PO<sub>4</sub> in the 8 x 10<sup>2</sup> - 3, 86 x 10<sup>10</sup> cps frequency range

SOURCE: Kristallografiya, v. 10, no. 5, 1965, 764-766

## TOPIC TAGS: rubidium compound, dielectric constant

ABSTRACT: The article reports on measurements of the complex dielectric constant  $\mathbb{E} * = \mathbb{E} + j\mathbb{E}'$ , i.e., measurements of the quantities  $\mathbb{E}_3 *$  and  $\mathbb{E}_1 * = \mathbb{E}_2 *$ , in RbH<sub>2</sub>PO<sub>4</sub> single crystals. Measurements at low and high frequencies were made with a low-frequency RFT bridge and a PIMEL low-capacitance meter, respectively. Measurements at 15 Me were made with a Ye9-5 Q-meter. It was found that the values of the electric constants of samples cut out of different portions of the same crystal are the same. However,  $\mathbb{E}'$  may change with the conditions of growth of the crystal. A marked decrease (dispersion) of the dielectric constant occurs in the 9.6 x 109 - 3.86 x 10<sup>10</sup> cps range; at these frequencies, a substantial increase in the loss-angle tangent begins. The nature of this dis-

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1011年日本地南北北京大学(1011年)古老地域市民社

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ACCESSION NR: AP5024569

persion can probably be determined by measurements at higher frequencies. Orig. art. has: 1 figure, 1 table, and 4 formulas.

ASSOCIATION: Institut kristallografii AN SSSR (Institute of Crystallography, AN SSSR)

SUBMITTED: 05Mar65

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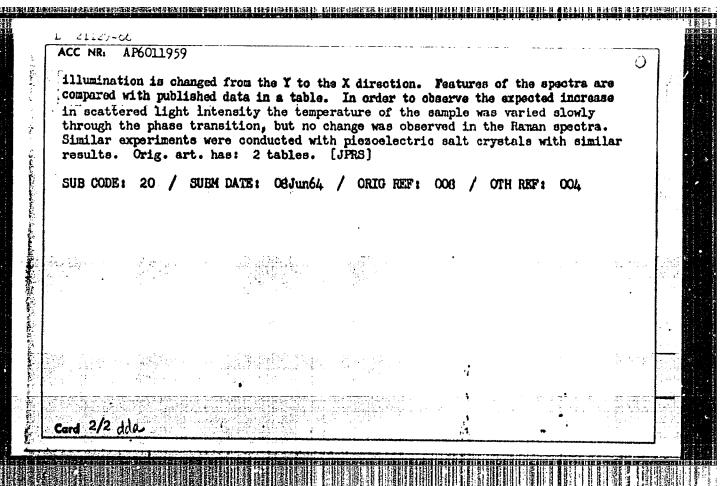
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L 21129-66 EWT( ACC NR. AP6011959 EWT(m)/EWP(j) RM SOURCE CODE: UR/0070/65/010/003/0335/0337 AUTHOR: Arbatskaya, A. N.; Zheludev, Y. S.; Zirnit, U. A.; Sushchinskiy, M. M. ORG: Institute of Crystallography, AN SSSR (Institut kristallografii AN SSSR) TITLE: Low-frequency vibrational spectra of triglycine sulphate and rochelle salt monocrystals during phase transitions SOURCE: Kristallografiya, v. 10, no. 3, 1965, 335-337 TOPIC TAGS: phase transition, Raman spectrum, Curie point, light scattering, crystal lattice vibration, quartz crystal, crystal symmetry, single crystal ABSTRACT: Raman scattering spectra of monocrystals of triglycine sulfate are studied near the Curie temperature (+490C); and those of Rochelle salt, near the upper Curie point (+24°C). According to the Ginsburg-Levanyuk theory, the spectral structure of the scattered light should change markedly near the phase transition points of the second kind. This change should be characterized by a decrease in the frequency of certain lattice vibrations when the temperature is raised. At the transition point of the second kind the frequency of these vibrations should become zero and the corresponding Raman lines intensify. Experiments with quartz appear to support the G-L theory. Triglycine sulfate undergoes a change in symmetry upon passing through the Curie point. A right-angle prismatic monocrystal was illuminated along the Y and then along the Z axis, and the scattered light was observed along the Z axis in both cases. A line at 47 cm appeared in the Raman spectra when the direction of UDC: 548.0: Card 1/2



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1,1 0,120% 2,1 0,120%	L 21222-66 EWT(m)/EWP(t) IJP(c) JD  ACC NR: AP6003812 SOURCE CODE: UR/0181/66/003/001/0272/0274	
	AUTHORS: Aref'vev, I. M.; Bazhulin, P. A.) (deceased); Gavrilova,	
	T. V.; Zheludev, I.S.	
	ORG: Physics Institute im. P. N. Lebedev AN SSSR, Moscow	
	(Fizicheskiy institut AN SSSR)	
	TITLE: Temperature dependence of the intensity of light scattering in oriented single crystals of KH2PO4 and Rochelle salt	
	SOURCE: Fizika tverdogo tela, v. 8, no. 1, 1966, 272-274	
	TOPIC TAGS: ferroelectric crystal, phase transition, light scattering, temperature dependence, light polarization, elastic modulus, crystal lattice vibration, Curie point, paraelectricity, piezoelectric	
	property	
	ABSTRACT: The purpose of the measurement of the temperature dependence was to check whether the ferroelectric phase transition in these ence was to check whether the ferroelectric phase transition in these ence was to check whether the ferroelectric phase transition in these ence was to check whether the ferroelectric phase transition in these ence was to check whether the ferroelectric phase transition in these ence was to check whether the ferroelectric phase transition in these ence was to check whether the ferroelectric phase transition in these ence was to check whether the ferroelectric phase transition in these ence was to check whether the ferroelectric phase transition in these ence was to check whether the ferroelectric phase transition in these ence was to check whether the ferroelectric phase transition in these ence was to check whether the ferroelectric phase transition in these ence was to check whether the ferroelectric phase transition in these ence was to check whether the ferroelectric phase transition in these ence was to check whether the ferroelectric phase transition in these ence was to check whether the ferroelectric phase transition in the ence was the ferroelectric phase transition in the ence was the ferroelectric phase transition in the ence was the ferroelectric phase transition.	
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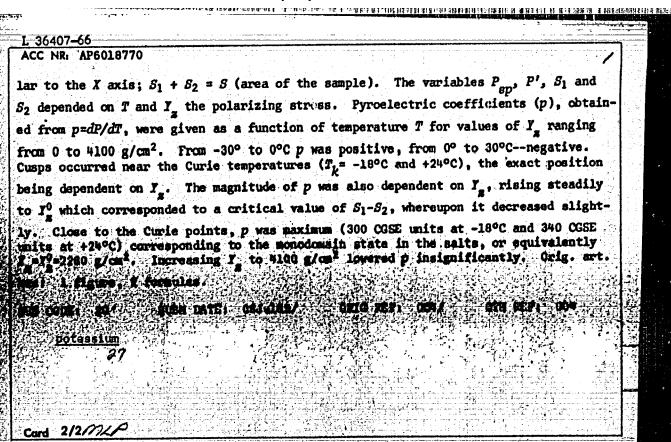
ACC NR: AP6003812

(DFS-12) whose output was photoelectrically recorded. tor and the sample-cooling system are described elsewhere (A. V. Rakov, Tr. FIAN v. 27, 111, 1964). The investigated crystals were transparent with cross sections 7.5 x 7.5 mm and lengths 20, 23, and The Rochelle-salt crystals measured 7.5 x 7.5 x 30 mm. The intensity of scattering was measured at the maximum of the Hg 4358 A line under smooth variation of the temperature. The results were strongly dependent on the polarization, and in the case of one type of polarization the intensity of the scattered light had a variation similar to that of the reciprocal of the elastic constant. It is concluded on this basis that the scattering is produced by anomalous acoustic vibrations. In the case of Rochelle salt, the effect is less pronounced in KH2PO4, and no increase in the scattering intensity is observed at the second Curie point. This indicates that the structure of the Rochelle salt crystal is different in the two paraelectric phases. No low-frequency Raman scattering spectrum was observed, and it is therefore deduced that the increase in the scattering intensity of the Curie point is connected with the anomalous behavior of the acoustic lattice vibrations. It is concluded on the

Card 2/3

L 21222-66 ACC NR: AP6003912
basis of these results and earlier data by the authors (PTT v. 7, 2413, 1965) that the ferroelectric phase transition in both salts is due to the instability of the crystal against the acoustic and optical vibrations of the lattice, which are interelated by the piezoeffect. The authors thank G. P. Motulevich and D. G. Sannikov for a useful discussion. Orig. art. has: 2 figures
SUB CODE: 20/ SUBM DATE: 02Aug65/ ORIG REF: 005/ OTH REF: 003
Card 3/34:

UR/0070/66/011/003/0415/0418 L 36407-66 SOURCE CODE: ACC NR. AP6018770 AUTHOR: Zheludey, I. S.; Gladkiy, V. V. ORG: Institute of Crystallography, AN SSSR (Institut kristallografii AN SSSR) of Rochelle salt The pyroelectric effect in single crystals SQURCE: Kristallografiya, v. 11, no. 3, 1966, 415-418 TOPIC TAGS: Rochelle salt, single crystal, pyroelectric effect, Curie temperature, elastic stress, polarisation ABSTRACT: Pyroelectric coefficients of Rochelle salt single crystals were measured between -30° and +30°C after polarization by mechanical straining. The pyroelectric charge was measured in a Tepler thermostat by the compensation method using slow heating rates. Temperatures were measured with Cu-constantan thermocouples to 1 or 2°C beyoud the Curie point. The polarization P of the strained crystals was obtained from  $P = P_{gp} \frac{S_1 - S_2}{S} + P'$ where  $P_{sp}$  is the spontaneous polarization; P' is the elastic polarization;  $S_1$ ,  $S_2$  are cross-sectional areas of domains having unlike signs in the crystal face, perpendicu-UDC: 548.0 : 537.227 Card



1 2437 -66 EWT(m)/EWA(d)/T/EWP(t) | 13P(c) ACC NR: AP6010980 SOURCE CODE: UR/0056/66/050/003/0595/0604 AUTHORS: Yamzin, I. I.; Stzov, R. A.; Zheludev, I. S.: Perekalina, T. M.; Zalesskiy, A. V. ORG: Institute of Crystallography, Academy of Sciences SSSR (Institut kristallografii Akademii nauk SSSR) TITLE: Spin ordering and magnetocrystalline anisotropy in single crystals of BaCoxFe18-xC27 ferrites SOURCE: Zhurnal eksperimentalinov i teoreticheskov fiziki, v. 50, no. 3, 1966, 595-604 TOPIC TAGS: ferrite, single crystal, magnetic anisotropy, neutron diffraction, nuclear spin, Curie point, temperature dependence, spin wave theory ABSTRACT: This is a continuation of earlier work by the authors (ZhETF v. 46, 1985, 1964). In this paper new data are presented on the magnetic anisotropy energy of the ferrite system under discussion. The crystals were grown by the Verneuil method and were the same as Card

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1 24777-56

used in the earlier investigation. In view of the fact that the fer rites investigated exhibit various types of magnetic anisotropy at low temperatures, the authors used a neutron diffraction method to investigate the influence of the cobalt ions on the positions of the spin ordering axis in these crystals in the temperature range from 77K to the Curie temperature. The temperature dependence of the magnetic anisotropy constants was investigated in the same range of temperatures and compared with the theory. The same samples were used to obtain neutron diffraction patterns as were used in the investigation of the magnetic anisotropy. The results show that the spin directions coincide with the directions of the total magnetization vectors of the crystals. The data also indicate that the experimental results can be fully reconciled with a theoretical formula deduced by Ye. A. Turov from the phenomenological theory of spin waves (Fizicheskiye svoystva magnitouporyadochennykh kristallov [Physical Properties of Mangetically Ordered Crystals], AN SSSR, 1963), without need to make allowance for any particular structure model. Orig. art. has: 7 figures, 3 formulas, and 3 tables.

SUB CODE: 20/ SUBM DATE: 250ct65/ ORIO REF: 003/ OTH REF: 009

Card 2/2 00 R

ENT(1)/ENT(m)/ENP(n)/EEC(k)-2/T/ENP(t)ACC NR AP6007268 SOURCE CODE: UR/0053/66/088/002/0253/0286 AUTHOR: Zheludev, I. 8. ORG: Institute of Crystallography, AN SSSR (Institut kristallografii AN SSSR) TITLE: Electro-ordical phenomena in crystals Ē SOURCE: Uspekhi fizicheskikh nauk, v. 88, no. 2, 1966, 253-286 TOPIC TAGS: electrooptical effect, refractive index, absorption coefficient, light absorption, electric field, cryatal ABSTRACT: This is a review article dealing with electro-optical effects, defined as the changes induced in the optical properties (refractive index and absorption coefficient) by application of an electric field. The different types of electro-optical effects are described (linear, quadratic, spontaneous), and the equations for these effects are derived in a matrix formulation for crystals of various symmetries (orthorhombic, monoclinic, triclinic, tetragonal, trigonal, and hexagonal). This is followed by a discussion of the optical systems used for the investigation of the electro-optical effect, and methods used to determine directly the electro-optical coefficients. Data are presented on the linear electro-optical effect in linear dielectrics and ferroelectrics, and on the quadratic electro-optical effect in crystals with central symmetry. Practical uses of the electro-optical properties of crystals for high-frequency light modulation and optical shutters, for sound recording, and for narrow-band interference-polarization light filters are discussed. Original art. has: 16 figures, 36 formulas, and 5 tables. SUB CODE: ODE: 20/ ORIG RED ORIG REF: 011/ OTH REF: 057 UDC: 535.5

L 09461-67 ACC NR: AP6024666

SOURCE CODE: UR/0070/66/011/004/0604/0609

AUTHOR: Gorelik, V.S.; Zheludev, I. S.; Sushchinskiy, M. M.

ORIG: Physics Institute im. P. N. Lebedev AN SSSR (Fizicheskiy institut AN SSSR); Institute of Crystallography AN SSSR (Institut kristallografii AN SSSR)

TITLE: Study of the Raman spectrum of NaNO2 single crystal near the phase transition point.

SOURCE: Kristallografiya, v. 11, no. 4, 1966, 604-609

TOPIC TAGS: sodium compound, Raman spectrum, phase transition, ferroelectricity, temperature dependence, line broadening, crystal lattice vibration

ABSTRACT: This is a continuation of an earlier study (Kristallografiya v. 10, no. 3, 335, 1965) and deals with the behavior of most lines of the Raman spectrum of single-crystal NaNO2 in the temperature interval from 30 to 178C. Principal attention was paid to the small temperature range (± 20°) near the phase transition point (160C). The single crystal was a rectangular prism 3 x 7 x 10 mm, cut so that its smallest side was oriented along the ferroelectric axis z. The Raman appetra were photographed with a spectrograph, using the 4358 Å mercury line for excitation.

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UDC: 548.0:535.36

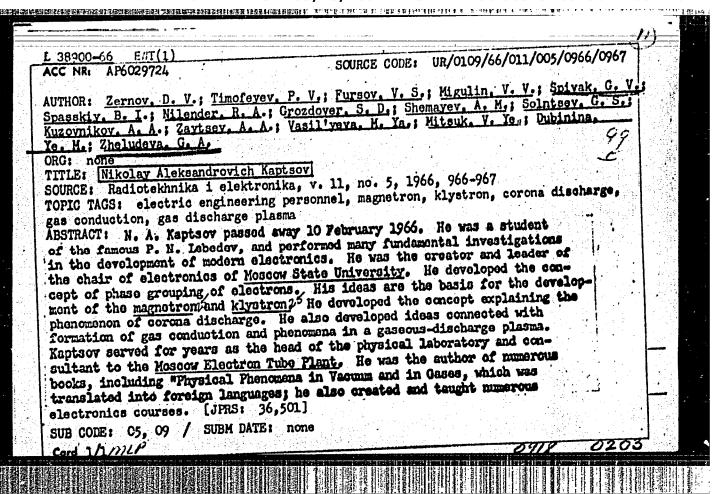
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L 09461-67 ACC NR: AP6024666

Soven lines were registered, whose frequencies agreed essentially with those published earlier. The Raman spectrum obtained near the transition point differed noticeably from that obtained at 30C. The low-frequency lines shifted in linear fashion, while the higher frequency lines exhibited practically no shift. All observed lines broadened with increasing temperature, but the broadening of the low-frequency lines was larger. Some of the lines vanished with increasing temperature. A group-theoretical analysis of the spectrum for both the high and the low frequency parts of the spectrum is used to interpret the results. The vanishing and the intensity variations of the spectra agree with the selection rules, and the broadening is due to ordinary temperature effects connected with the increase of the interaction between the lattice oscillators themselves and the interaction between the lattice oscillators and other degrees of freedom of the crystal. The authors thank Professor P. A. Bazhulin for interest in the work and also V. I. Mursin for valuable advice. Orig. art. has: 4 figures, 1 formula, and 2 tables.

SUB CODE: 20/ SUEM DATE: 03 AUG 65 ORIG REF: 012/ OTH REF: 005

Card 2/2 1



L 09462-67 EWT(1)/EWT(m)/EWP(t)/ETI IJP(c) JD/GG ACC NR: APG024667 SOURCE CODE: UR/0070/66/011/004/0610	
AUTIOR: Zheludov, I. S.; Romanyuk, N. A.	<b>4</b>
ORIG: L'vov Stato University (L'vovskiy gosudarstvennyy universitet); In Crystallography AN SSSR (Institut kristallografii AN SSSR)	stitute of,
TITLE: Diolectric properties of clamped Rochello-salt crystals	
SOURCE: Kristallografiya, v. 11, no. 4, 1966, 610-613	
TOPIC TAGS: forroclectric crystal, piozoclectric crystal, electric hyst diolectric constant, pressure effect, Curie point, electric polarization	orosis,
ABSTRACT: The authors report an investigation of the domain structure, thysteresis loop, and the initial dielectric constant of Rochelle-salt critical region of the upper Curie temperature for different degrees of comparts the samples. The tests were made on sugare polished plates with 450 X-cuing $10 \times 10 \times (0.4 0.7)$ mm. The samples were measured on a microscope a thermostat equipped with a device for simultaneous or successive comparts the crystals in two mutually perpendicular directions. This apparatus region in an earlier paper (Kristallografiya v. 4, no. 5, 710 717, 19	rossion of nt, moasur- n stage in rossion of was des-

L 09462-67 ACC NR: AP6024667

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dielectric constant was measured with a bridge and a hysteresis loop was by means of a Sawyer-Tower circuit. The domain structure was observed visually and photographed when necessary. The results showed that the dielectric constant decreased with increasing compression, the upper Curie point increased in the case of unilateral compression and decreased in the case of bilateral compression, and the hysteresis loops gradually contracted to lines with increasing compression. The results agree with modern theoretical notions concerning the laws governing polarization of ferroelectrics and explain the radiative changes occurring in the dielectric properties of ferroelectric materials. The results also serve as a confirmation of the frequently used "internal field" model for the explanation of the polarization of the ferroelectric. Orig. art. has: 3 figures.

SUB CODE: 20/

SUBM DATE: 15Jul65/

ORIG REF: 003/

OTH REF: 006

Card 2/2-2C

L 09461-67

ACC NR: AP6024656

SOURCE CODE: UR/0070/66/011/004/0604/0609

AUTHOR: Gorolik, V.S.; Zholudev, I. S.; Sushchinskiy, M. M.

ORIG: Physics Institute im. P. N. Lobedov AN SSSR (Fizicheskiy institut AN SSSR) Institute of Crystallography AN SSSR (Institut kristallografii AN SSSR)

TITLE: Study of the Raman spectrum of NaNO2 single crystal near the phase transition point.

SOURCE: Kristallografiya, v. 11, no. 4, 1966, 604-609

TOPIC TAGS: sodium compound, Raman spectrum, phase transition, ferroelectricity, temperature dependence, line broadening, crystal lattice vibration

ABSTRACT: This is a continuation of an earlier study (Kristallografiya v. 10, no. 3, 335, 1965) and deals with the behavior of most lines of the Raman spectrum of single-crystal NaNO2 in the temperature interval from 30 to 178C. Principal attention was paid to the small temperature range (± 200) near the phase transition point (160C). The single crystal was a rectangular prism 3 x 7 x 10 mm, cut so that its smallest side was oriented along the ferroelectric axis s. The Raman spectra were photographed with a spectrograph, using the 4358 Å mercury line for excitation.

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VDC: 548.0:535.36

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L 09461-67

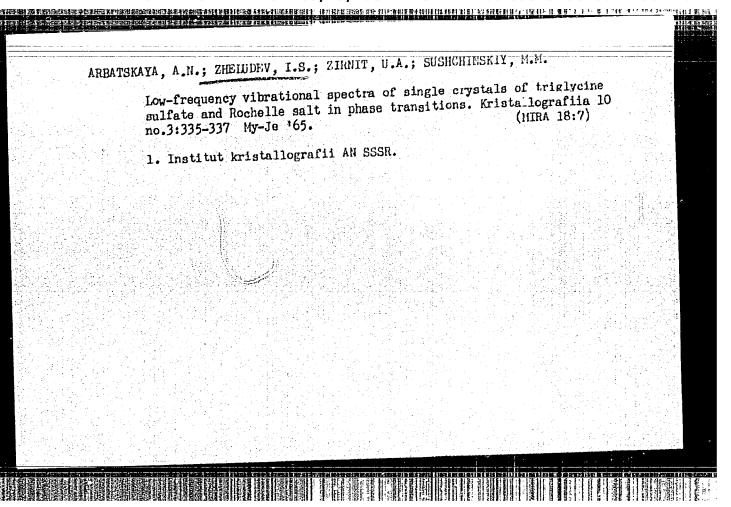
ACC NR: AP6024666

Soven lines were registered, whose frequencies agreed essentially with these published earlier. The Raman spectrum obtained near the transition point differed noticeably from that obtained at 30°C. The low-frequency lines shifted in linear fashion, while the higher frequency lines exhibited practically no shift. All observed lines broadened with increasing temperature, but the broadening of the low-frequency lines was larger. Some of the lines vanished with increasing temperature. A group-theoretical analysis of the spectrum for both the high and the low frequency parts of the spectrum is used to interpret the results. The vanishing and the intensity variations of the spectra agree with the selection rules, and the broadening is due to ordinary temperature effects connected with the increase of the interaction between the lattice oscillators themselves and the interaction between the lattice oscillators and other degrees of freedom of the crystal. The authors thank Professor P. A. Bazhulin for interest in the work and also V. I. Mursin for valuable advice. Orig. art. has: 4 figures, 1 formula, and 2 tables.

SUB CODE: 20/ SUBM DATE: 03 AUG 65 ORIG REF: 012/ OTH REF: 005

Card 2/2 1C

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	Electro-optical effects in crystals. Usp. fiz. nauk 88 no.2: 253-286 F '66. (MIRA 19:2)	
	l. Institut kristallografii AN SSSR.	



ZAYTSEVA, M.P.; ZHELIJDEV; I.S.; ZHEREBTSOVA; L.I.; NOICHENKOV, A.A.

Intensity of an electric field required to bring about polarization equal to spontaneous polarization. Izv. AN SSSR. Ser. fiz. 29 nc.6:948-950 Je '65. (MIRA 18:6)

1. Institut fiziki Sibirskogo otdeleniya AN SSSR i Institut kristallografii AN SSSR.

12 F.	的现在分词,我们就是我们的现在分词,我们就是一个人的,我们就是一个人的,我们就是一个人的,我们就是一个人的,我们就是一个人的,我们就是一个人的,我们就是一个人的 第二章
	L. 7986-66 Edit(d)/Edit(1), Edit(k)-2  AUTHORS: Vlokh, O. G.; Zheludev, I. S.; Shemburov, V. A.
	ORG: none
	TITLE: Electrooptical modulator. Q Class 42, No. 175272
:	SOURCE: Byulleten' izobreteniy i tovarnykh znakov, no. 19, 1965, 85
	0. 44 5
	TOPIC TAGS: electrooptic effect, electric field
:	ABSTRACT: This Author Certificate describes an electrooptical modulator consisting of crossed polarizers between which is situated a crystal in an electric field. The direction of the electric field is parallel to the direction of light and the axis of symmetry. To eliminate the treatment of the crystal surface and the influence of temperature and moisture of the surrounding medium on its performance and also to eliminate turning the optical axis through an angle of 22.5° under nonresonance condition, use is made of a pentaerythritol crystal.
7.	SUB CODE: OP/ SUBM DATE: 26Jan63
	Card 1/1 OC: 621.376.9
Щ	

AREF'YEV, I.M.; BAZHULIN, P.A.; ZHEIUDEV, I.S.

Long-wave infrared transmission spectra of NH, H<sub>2</sub> PO<sub>2</sub>. Fiz. tver. tela 7 no.9:2834-2836 S '65.

(MIRA 18:10)

1. Pizicheskiy institut imeni P.N.Lebedeva AN SSSR, Hoskva.

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<b>Z</b> H)	BLUDEV, N.A., kandidat tekhnicheskikh nauk.	E M ( 55 )
	Computation of corrections for the calculations of natural pitching and rolling of ships at anchor. Shor.trud.Len.Gidroneteorol.inst. no.2:93-120 (MLRA 6:8)  (Anchorage) (Stability of ships)	
		5

#### ZHELUDEV, N.A.

"Determination of the Equilibrium Form of a Cable Loaded with n loads" Sb. Tr. Leningr. Gidrometeorol. In-ta, No. 3, 1954, i25-147

The author carries out the theoretical determination of the equilibrium form of a cable attached at one end with loads suspended alon; its length. The cable is in a steady state horizontal homogenous flow. The Equilibrium equations are integrated under various assumptions. The author derives formulas for the approximate determination of the tensile strength and equilibrium form of the cable. (RZhMekh, no. 9, 1955)

sov/179-59-3-38/45 Zheludev, P. I. (Moscow) AUTHOR: On the Supersonic Flow Past a Flat Quasi-triangular TITLE: Short Wing (K sverkhzvukovomu obtekaniyu ploskikh kvazitreugol'nykh kryl'yev malogo udlineniya) PERIODICAL: Izvestiya Akademii nauk SSSR, Otdeleniye tekhnicheskikh nauk, Mekhanika i mashinostroyeniye, 1959, Nr 3, pp 202-203 (USSR) ABSTRACT: The work is based on Zheludev's work (Ref 1) describing a method of determining the velocity potential in supersonic flow past thin bodies. The complex potential, Eq (1), is taken from Ref 2, where 2a(x) - wing span (- and + = upper and lower half-planes respectively). The limiting conditions for the velocity potential is given as Eq (2), the partial solution of which is defined as Eqs (3) and (4) and their harmonic functions as Eqs (5) and (6). The latter can be expressed in the form of the series (7) and (8). The coefficient of pressure for the triangular wing, c, can be shown as Eq (9) for  $a(x) = x/tg \chi$ , where X - the wing's Card 1/2 angle, z - coordinates of the span.

**APPROVED FOR RELEASE: 03/15/2001** CIA-RDP86-00513R002064710012-8"

This coefficient

SOV/179-59-3-38/45
On the Supersonic Flow Past a Flat Quasi-triangular Short Wing is given as Eq (10) in Ref 1, where  $z \neq a(x)$ . There are 2 references, 1 of which is Soviet and 1 English.

SUBMITTED: December 10, 1958

Card 2/2

AUTHOR: Zneludev, P. I. (Moscow)		6.7	
ORG: none TITLE: Aerodynamic characteristics of a thin finned body	ाक्ताम् 🖟	57 B	
SOURCE: AN SSSR. Izvestiya. Mekhanika, no. 6, 1965, 122-123			
TOPIC TAGS: aerodynamics, aerodynamic characteristic, aerodyna plex function	mic coefficient,	COM-	
ABSTRACT: The aerodynamic moments and coefficients are calcula body at angle of attack and side-slip in supersonic flow (see F	tied for a cross-	finned	
Fig. 1.		er pr	
Card 1/2			

12

A potential flow analysis is used with the flow potential defined by

$$w(\zeta) = \frac{1}{2\pi} \frac{a\beta}{dz} \ln (\zeta - \zeta_0) + \alpha \left[ \left( \zeta - \zeta_0 - \frac{R^4}{\zeta - \zeta_0} \right)^2 + \left( a + \frac{R^4}{a} \right)^2 \right]^{1/4} + i\beta \left[ \left( \zeta - \zeta_0 - \frac{R^4}{\zeta - \zeta_0} \right)^2 - \left( z - \frac{R^4}{\zeta - \zeta_0} \right)^2 \right]^{1/4} + i\beta \left[ \left( \zeta - \zeta_0 - \frac{R^4}{\zeta - \zeta_0} \right)^2 \right]^{1/4} + i\beta \left[ \left( \zeta - \zeta_0 - \frac{R^4}{\zeta - \zeta_0} \right)^2 \right]^{1/4} + i\beta \left[ \left( \zeta - \zeta_0 - \frac{R^4}{\zeta - \zeta_0} \right)^2 \right]^{1/4} + i\beta \left[ \left( \zeta - \zeta_0 - \frac{R^4}{\zeta - \zeta_0} \right)^2 \right]^{1/4} + i\beta \left[ \left( \zeta - \zeta_0 - \frac{R^4}{\zeta - \zeta_0} \right)^2 \right]^{1/4} + i\beta \left[ \left( \zeta - \zeta_0 - \frac{R^4}{\zeta - \zeta_0} \right)^2 \right]^{1/4} + i\beta \left[ \left( \zeta - \zeta_0 - \frac{R^4}{\zeta - \zeta_0} \right)^2 \right]^{1/4} + i\beta \left[ \left( \zeta - \zeta_0 - \frac{R^4}{\zeta - \zeta_0} \right)^2 \right]^{1/4} + i\beta \left[ \left( \zeta - \zeta_0 - \frac{R^4}{\zeta - \zeta_0} \right)^2 \right]^{1/4} + i\beta \left[ \left( \zeta - \zeta_0 - \frac{R^4}{\zeta - \zeta_0} \right)^2 \right]^{1/4} + i\beta \left[ \left( \zeta - \zeta_0 - \frac{R^4}{\zeta - \zeta_0} \right)^2 \right]^{1/4} + i\beta \left[ \left( \zeta - \zeta_0 - \frac{R^4}{\zeta - \zeta_0} \right)^2 \right]^{1/4} + i\beta \left[ \left( \zeta - \zeta_0 - \frac{R^4}{\zeta - \zeta_0} \right)^2 \right]^{1/4} + i\beta \left[ \left( \zeta - \zeta_0 - \frac{R^4}{\zeta - \zeta_0} \right)^2 \right]^{1/4} + i\beta \left[ \left( \zeta - \zeta_0 - \frac{R^4}{\zeta - \zeta_0} \right)^2 \right]^{1/4} + i\beta \left[ \left( \zeta - \zeta_0 - \frac{R^4}{\zeta - \zeta_0} \right)^2 \right]^{1/4} + i\beta \left[ \left( \zeta - \zeta_0 - \frac{R^4}{\zeta - \zeta_0} \right)^2 \right]^{1/4} + i\beta \left[ \left( \zeta - \zeta_0 - \frac{R^4}{\zeta - \zeta_0} \right)^2 \right]^{1/4} + i\beta \left[ \left( \zeta - \zeta_0 - \frac{R^4}{\zeta - \zeta_0} \right)^2 \right]^{1/4} + i\beta \left[ \left( \zeta - \zeta_0 - \frac{R^4}{\zeta - \zeta_0} \right)^2 \right]^{1/4} + i\beta \left[ \left( \zeta - \zeta_0 - \frac{R^4}{\zeta - \zeta_0} \right)^2 \right]^{1/4} + i\beta \left[ \left( \zeta - \zeta_0 - \frac{R^4}{\zeta - \zeta_0} \right)^2 \right]^{1/4} + i\beta \left[ \left( \zeta - \zeta_0 - \frac{R^4}{\zeta - \zeta_0} \right)^2 \right]^{1/4} + i\beta \left[ \left( \zeta - \zeta_0 - \frac{R^4}{\zeta - \zeta_0} \right)^2 \right]^{1/4} + i\beta \left[ \left( \zeta - \zeta_0 - \frac{R^4}{\zeta - \zeta_0} \right)^2 \right]^{1/4} + i\beta \left[ \left( \zeta - \zeta_0 - \frac{R^4}{\zeta - \zeta_0} \right)^2 \right]^{1/4} + i\beta \left[ \left( \zeta - \zeta_0 - \frac{R^4}{\zeta - \zeta_0} \right)^2 \right]^{1/4} + i\beta \left[ \left( \zeta - \zeta_0 - \frac{R^4}{\zeta - \zeta_0} \right)^2 \right]^{1/4} + i\beta \left[ \left( \zeta - \zeta_0 - \frac{R^4}{\zeta - \zeta_0} \right)^2 \right]^{1/4} + i\beta \left[ \left( \zeta - \zeta_0 - \frac{R^4}{\zeta - \zeta_0} \right)^2 \right]^{1/4} + i\beta \left[ \left( \zeta - \zeta_0 - \frac{R^4}{\zeta - \zeta_0} \right)^2 \right]^{1/4} + i\beta \left[ \left( \zeta - \zeta_0 - \frac{R^4}{\zeta - \zeta_0} \right)^2 \right]^{1/4} + i\beta \left[ \left( \zeta - \zeta_0 - \frac{R^4}{\zeta - \zeta_0} \right)^2 \right]^{1/4} + i\beta \left[ \left( \zeta - \zeta_0 - \frac{R^4}{\zeta - \zeta_0} \right)^2 \right]^{1/4} + i\beta \left[ \left( \zeta - \zeta_0 - \frac{R^4}{\zeta - \zeta_0} \right)^2 \right]^{1/4} + i\beta \left[ \left( \zeta - \zeta_0 - \frac{R^4}{\zeta - \zeta_0} \right)^2 \right]^{1/4} + i\beta \left[ \left( \zeta -$$

This potential is then expanded in inverse powers of ( and the following forces and moments are calculated: the side and lift forces Z, Y, the three moments N, My, Mg, and the corresponding coefficients, e.g.,

$$C_{\nu} = \frac{2z \left(b^4 + R_1^4 - R_1^2 b^3\right)}{R_1^2 b^3}, \qquad C_{z} = \frac{23 \left(b^4 + R_1^4 + R_1^2 b^3\right)}{R_1^2 b^3}$$

$$C_{mx} = -4\alpha\beta \left[1 - \frac{1}{R_1^3} \int_{0}^{1} R^3(x) dx\right].$$

Orig. art. has: 8 equations and 2 figures.

SUB CODE: 20/

ACC NR. AP6002327

SUBM DATE: 10Jul64/

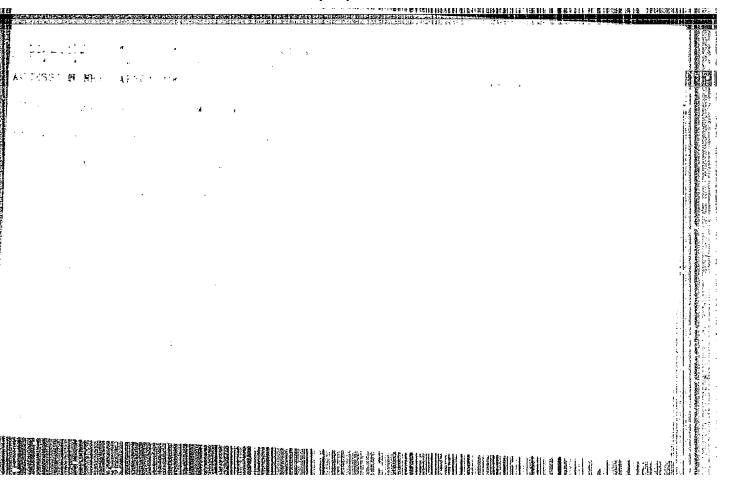
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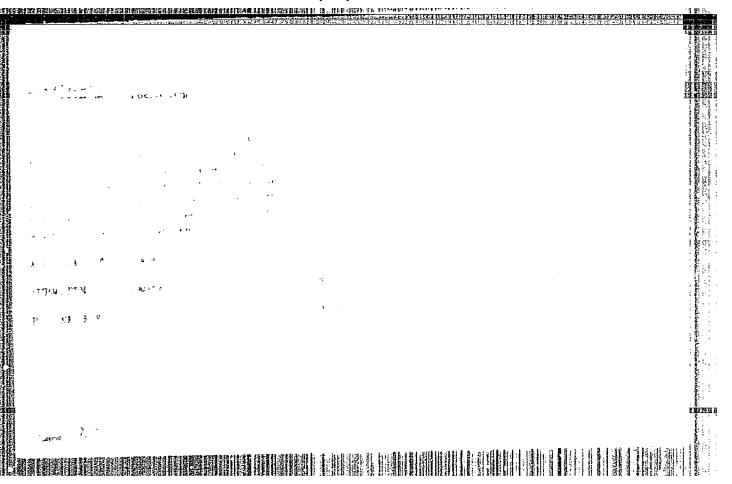
Card 2/2 FW

	Unsteady supersonic flow about bodies of revolution a bending flutter: Izv.AN SSSR. Mekh.i mashinostr. no Ja-F '64.	and their p.1:158-162 (MIRA 17:4)	

		<u> </u>	
CCESSION NR: AP4018437	B/0179/64/000/001/0168/0162	İ	ر
UTHOR: Zheludev, P. I. (Moscow)			
ITLE: Unsteady flow past and flexing flutter conic stream	of solids of revolution in a super-		
OURCE: AN SSSR. Izv. Otd. tekh. nauk. Mekhan 964, 158-162	nika i mashinostroyeniye, no. 1,		
OPIC TAGS: fluid dynamics, fluid mechanics, as lexure	rodynamics, unsteady flow, flutter,		
BSTRACT: In the statement of the theory of a tonic Flow Past Slender Pointed Bodies, Quart. Go. 75) there is an examination of supersonic fluscribe harmonic elastic oscillations and expresential, pressure coefficient and the transverength of the solid. In this article there is a lexing oscillations of solids of revolution such free boundaries and beams with one tightly a given from an earlier determined expression f	. Mech. Appl. Math., 1949, vol. 2, ow past solids of revolution which ssions are found for the velocity se force distributed along the further examination of the elastic h as beams of variable cross section closed end. The acrodynamic force		
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		100	H Fig.

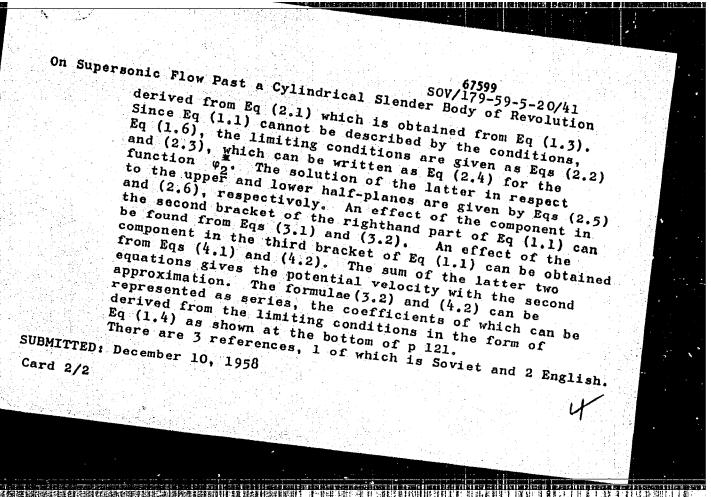
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solved by the Bubnov-Galerki: the natural oscillations and	The differential equation for the method. An expression is obtain examples are given of calculation nical beam. Orig. art. has: 2 fi	ed for calculation of of the speed of	
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ZHELUDEV, P. I., Candidate Tech Sci (diss) -- "Supersonic flow around certain thin bodies". Moscow, 1959. 5 pp (Acad Sci USSR, Dept of Tech Sci, Inst of Mech), 130 copies (KL, No 26, 1959, 125)

67599 sov/179-59-5-20/41 10,6000 P. I. (Moscow) AUTHOR: On Supersonic Flow Past a Cylindrical Slender Body of TITLE: Revolution PERIODICAL: Izvestiya Akademii nauk SSSR, Otdeleniye tekhnicheskikh nauk, Mekhanika i mashinostroyeniye, 1959, Nr 5, pp 118-121 (USSR) ABSTRACT: The flow past a thin pointed body of revolution with thin wings of small aspect ratio is considered. The equation of motion is given as Eq (1.1) for the limiting conditions of the potential velocity  $\phi$  =  $\phi$  +  $\phi$ 2 defined as Eq (1.2). After substituting the derivates of  $\varphi_0 = \frac{1}{2} (w + \bar{w})$  in the right term of Eq (1.1), the Poisson equation can be obtained, the solution of which can be shown as Eq (1.3). The potential velocity with the first approximation can be shown as Eq (1.5), where  $2a(\chi)$  - span of wing,  $\alpha$  - angle between the axis of the body and an undisturbed flow, R = const - radius of the body. The limiting conditions for  $\phi_2$  can be derived from Eq (1.2) and presented as Eq (1.5,a) for the surface of the body, and as Eq (1.6b) for the surface of the wing. An effect of the term  $M_{\infty}^2 - 1) \partial^2 \varphi / \partial \chi^2$  in Eq (1.1) can be Card 1/2



AUTHOR: Zheludev, P.I. (Moscow) 80V/24-58-9-10/31

TITIE: The Supersonic Flow Round Slender Bodies of Rotation With and Without Tail Units (Sverkhzvukovoye obtekaniye operennykh i neoperennykh tonkikh tel vrashcheniya)

PERIODICAL: Izvestiya Akademii Nauk SSSR, Otdeleniye Tekhnicheskikh Nauk, 1958, Nr 9, pp 74 - 82 (USSR)

ABSTRACT: The bodies are assumed to have an angle of attack α and to have their tail units inclined at an angle β to the vertical. A solution for the second approximation in slender body theory (Ref 1) is constructed. The frame of reference is chosen with the y axis vertically upwards and the velocity in the direction of increasing x. Forces on bodies of revolution having 2, 3 or 4 plane wings of small thickness with their trailing edges in the plane x = 1 are given. The solution from which these are derived is regarded as the first approximation and it is shown that the problem of finding the second approximation can be reduced to that of finding harmonic functions (as in the case of the first approximation). Hence, it is possible to solve the problem of finding the velocity potential of the second approximation for bodies of arbitrary transverse and longitudinal cross-

The Supersonic Flow Round Slender Bodies of Rotation With and Without Tail Units

sections satisfying conditions imposed on the maximum thickness and angle which any tangent plane to the surface of the body makes with the direction of the undisturbed flow. These quantities must be of the order of the maximum thickness of the body which in its turn must be very much less than the length of the body. In addition, the radius of curvature of an arbitrary transverse section must be of the order of  $\mathbf{d}^{-1}$ , where  $\mathbf{d}$  is the maximum thickness of the section. This condition is imposed on convex sections but not on concave ones. As an example, a circular cone of apex angle  $2\beta_1$  is considered. The expression for the pressure coefficient at zero angle of attack coincides with that obtained by Broderick (Ref 2) with the exception that a term of the form:

$$[(\chi + 1)/(M_{\infty}^2 - 1)]M_{\infty}^4$$

does not appear. This term can be introduced into the Card2/3 expression by using results of Van Dayk (Van Dyke) who

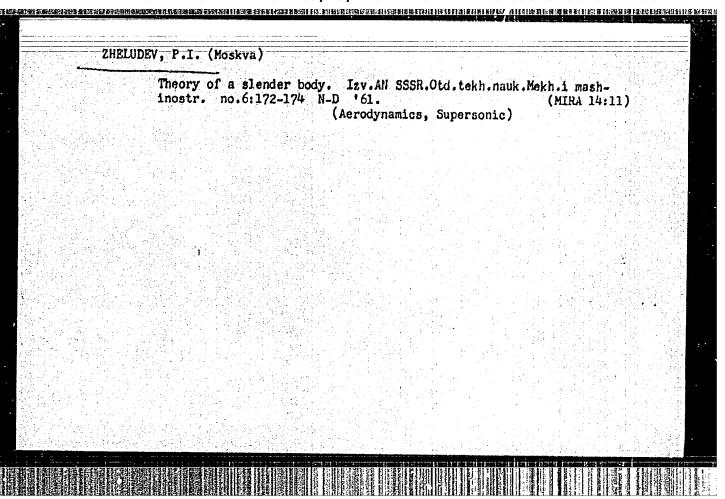
The Supersonic Flow Round Slender Bodies of Rotation With and Without Tail Units

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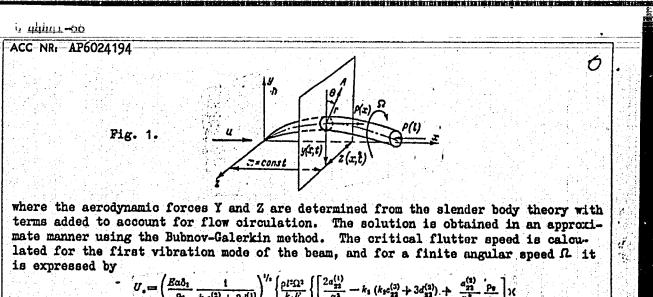
considered the case of an elliptic cone at zero angle of attack. The results for the revised formula are given graphically in which they are compared with the results of exact theory and the first approximation. The comparison shows that for sufficiently slender cones the second approximation gives better results than the first approximation. A similar graphical comparison is made where the abscissa is an angular co-ordinate in the plane x= constant rather than the Mach number. The second approximation again gives good results and the overall conclusion is that for  $\beta_1 < 15^\circ$  a better agreement with experiment can be expected from the second approximation.

There are 6 figures and 6 English references. SUBMITTED: February 7, 1958

Card 3/3



<u> L հիկի</u>1-66 EWT(1)/EWP(m)ACC NR. AP6024194 UR/0424/66/000/002/0160/0165 AUTHOR: Zheludev, P. I. (Moscow) ORG: none TITLE: Bending flutter of rotating elongating bodies Inzhenernyy zhurnal. Mekhanika tverdogo tela, no. 2, 1966, 160-165 SOURCE: TOPIC TAGS: elasticity theory, supersonic flow, unsteady flow, aeroelasticity, vibration ABSTRACT: An analysis is made to determine the effect of longitudinal acceleration and spinning motion on the elastic bending vibrations of an elongating body in a low supersonic flow. As shown in Fig. 1, small unsteady motions y(x,t) and z(x,t) are added to the variable longitudinal motion U and the body is assumed to be an elastic beam of variable cross section. The governing equations of motion are given by  $\frac{\partial^{3}}{\partial x^{3}} \left[ EI(x) \frac{\partial^{2} y}{\partial x^{3}} \right] + \frac{\partial}{\partial x} \left[ P(x) \frac{\partial y}{\partial x} \right] + \rho S(x) \frac{\partial^{2} y}{\partial t^{3}} - \frac{\partial}{\partial x} \left[ \rho I(x) \frac{\partial^{3} y}{\partial x \partial t^{3}} \right] - \frac{\Omega^{2} l}{\Omega x^{3}} \left[ \rho I(x) \frac{\partial y}{\partial x} \right] - 2\Omega \rho S(x) \frac{\partial z}{\partial t} - Y(x, t)$   $\frac{\partial^{3}}{\partial z^{3}} \left[ EI(x) \frac{\partial^{2} z}{\partial x^{3}} \right] + \frac{\partial}{\partial x} \left[ P(x) \frac{\partial z}{\partial x} \right] + \rho S(x) \frac{\partial^{3} z}{\partial t^{3}} - \frac{\partial}{\partial x} \left[ \rho I(x) \frac{\partial^{3} z}{\partial x \partial t^{3}} \right] - \frac{\Omega^{3} l}{\Omega x^{3}} \left[ \rho I(x) \frac{\partial z}{\partial x} \right] + 2\Omega \rho S(x) \frac{\partial y}{\partial t} - Z_{s}(x, t)$ Card 1/2



$$\begin{split} U_{\bullet} &= \left(\frac{E\alpha\delta_{1}}{\rho_{0}} \frac{1}{k_{2}c_{23}^{(2)} + 2d_{12}^{(1)}}\right)^{1/s} \left\{\frac{\rho l^{2}\Omega^{2}}{k_{2}E} \left\{ \left[\frac{2a_{23}^{(1)}}{\alpha^{3}} - k_{1}\left(k_{2}c_{22}^{(2)} + 3d_{33}^{(2)}\right) - \frac{a_{23}^{(3)}}{\alpha\delta_{1}} \frac{\rho_{9}}{\rho}\right] \right\} \\ &\times \left(\frac{k_{2}d_{22}^{(2)}}{k_{2}d_{23}^{(2)} + a_{22}^{(1)}}\right)^{3} + \frac{2}{\alpha} \left(\frac{2a_{23}^{(1)}}{\alpha} - \frac{a_{23}^{(2)}}{\delta_{1}} \frac{\rho_{9}}{\rho}\right) \frac{k_{1}d_{33}^{(2)}}{k_{2}d_{23}^{(2)} + a_{23}^{(1)}} - \frac{a_{23}^{(2)}}{k_{2}d_{23}^{(2)} + \delta_{23}^{(2)}}\right\}^{1/s} \\ &- k_{2}\left[k_{3}\left(k_{2}\delta_{22}^{(1)} + \delta_{22}^{(2)}\right) + \delta_{22}^{(2)}\right] - k_{2}\left[k_{2}\left(k_{2}a_{23}^{(2)} + \delta_{23}^{(2)}\right) + \delta_{23}^{(1)}\right]^{1/s} \\ &- k_{3}\left[k_{3}\left(k_{2}\delta_{23}^{(1)} + \delta_{23}^{(2)}\right) + \delta_{23}^{(2)}\right] + \frac{a_{23}^{(1)}}{\alpha} - \frac{a_{23}^{(1)}}{\alpha} - \frac{a_{23}^{(2)}}{\beta} - \frac{a_{2$$

Orig. art. has: 23 equations and 1 figure.

SUB CODE: 20/ SUBM DATE: 15Nov65/ ORIG REF: 002/ OTH REF: 002

APPROVED FOR RELEASE: 03/15/2001

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FILIPENKO, A.N. (st.Torshok, Oktyabr'skoy dorogi); ZHELUDEV, R.I. (st.Torshok, Oktyabr'skoy dorogi)

Chemical weed control on the track. Put'i put.khoz. 6 no.6:25-26 '62.

(MIRA 15:7)

(Railroads--Maintenance and repair) (Herbicides)

KOLESOV, A.P., prof.; ZHELIDEV, S.I.; DAVIDENKO, V.A.

Mediastinal and mediastinal-pulmonary form of sarcoidosis in the surgical clinic. Khirurgita 40 no.1:11-16 Ja '64.

(MIRA 17:11)

1. Khirurgitheskaya klinika dlya unovershenstvovaniya vrachey
No.1 Voyenno-meditsinskoy ordena Lenina akademii imeni Kirova.

ZHELUDEVA, G.A.

AUTHOR: Zheludeva, G.A.

109-3-10/23

TITIE: Work Function of the Antimony-caesium Cathode (O rabote, vykhoda sur'myano-tseziyevogo katoda)

PERIODICAL: Radiotekhnika i Elektronika, 1958, Vol. III, No. 3, pp. 395 - 399 (USSR).

ABSTRACT: The investigation reported was primarily concerned with the determination of the photo-electric work function of the Sb-Cs cathode as a function of the wavelength of the illuminating source. The magnitude of the work function was evaluated from the Einstein equation and attention was focused on the magnitude of the relative value of the work function rather than its absolute value. The investigation was done by the spherical-condenser method and two types of photo-cathodes were prepared. One of the photo-cathodes was a cylinder having 1 cm dia. and a length of 1 cm, while the second cathode was in the form of a glass sphere of 1 cm diameter. The cathodes had the usual "dark cherry" colour if deposited on glass and a grey-green colour if deposited on nickel. The collectors of the spherical condenser were in the form of spherical bulbs, having a diameter of 12 or 15 cm and were coated with a layer of aluminium plus a thin layer of soot (to eliminate the reverse Card1/3 currents). The photo currents were measured by the potentiometer

Work Function of the Antimony-caesium Cathode

109-3-10/23

method by using an electrometer, so that a sensitivity of  $10^{-13}$  A could be obtained. A quartz monochromator was used as a source of spectral lines. Current voltage curves were taken for the following wavelengths: 5791, 5461, 4047, 3651, 3130, 2804, 2482, 2200 Å; the curves are shown in the figure on p. 397. The work function can be determined from these curves by finding their intersection points with the horizontal (voltage) axis. It was found that the work functions of both the cathodes are independent of the illuminating wavelength for the range 5791 to 4047 Å. In the range 3130 to 2200 Å, the work function is also constant, but its value is somewhat lower than that for the range of the longer waves. An attempt was also made to determine the "red boundary" of the photo effect and it was found that the Fowler method could not be used for this purpose. The author expresses gratitude to Professor N.A. Kaptsov for directing this work. There are 1 figure, 1 table and 10 references, 6 of which are Russian and 4 English.

ASSOCIATION:

- Chair of Electronics of Moscow State University imeni M.V. Lomonosov (Kafedra elektroniki Moskovskogo universiteta im. M.V. Lomonosova)

Card2/3

Work Function of the Antimony-caesium Cathode 109-3-10/23

SUBMITTED: December 19, 1956

AVAILABLE: Library of Congress
Card 3/3

ZHELUDEVA, A.A.

AUTHORS:

Zheludeva, G.A. and Akhmatova, N.A.

109-3-11/23

TITIE:

Energy Distribution of the Photo-electrons in the Antimony-caesium Cathode during Various Stages of its Formation (Raspredeleniye fotoelektronov po energiyam dlya sur'myano-tseziyevogo katoda na raznykh stadiyakh yego formirovaniya)

PERIODICAL: Radiotekhnika i Elektronika, 1958, Vol. III, No. 3, pp. 400 - 404 (USSR).

ABSTRACT: The investigation of the energy distribution of the photoelectrons was done by the spherical-condenser method. The
investigated cathodes were prepared as follows: a layer of Sb
was deposited by evaporation on to a glass sphere. The layer
had a thickness of 1 500 Å and was not transparent. The activation process followed the standard technique, i.e. the Sb
layer was treated by Cs vapours at a temperature of 180 °C.
Three types of the activated film were prepared. The first
stage was characterised by a straw-yellow colouring of the
emissive surface; the sensitivity of the surface was 0.1 §
that of the normal Sb-Cs cathode. The second film, corresponding
to the second activation stage, had a light-red colouring and
its sensitivity was 0.6 that of the standard cathode. Finally,
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Energy Distribution of the Photo-electrons in the Antimony-caesium Cathode during Various Stages of its Formation

sensitivity (is sensitivity was taken as unity). Current voltage curves of the three cathodes were taken at the following wavelengths: 5 461, 4 047, 2 805, 2 482 and 2 300 Å. The resulting curves are shown in Figs. 1a, 6 and b. The curves were used to evaluate the work functions of the cathode by employing the Einstein equation and it was found that these were 1.78 ± 0.05, 1.56 ± 0.05 and 1.51 ± 0.05 eV for the first, second and third activation stages, respectively. These values are valid for the wavelength of 2 805 Å. By differentiating the current voltage curves of Fig. 1, it was possible to obtain the curves of the photo-electron energy distribution; these are shown in Figs. 2a, 6 and b; curves of Fig. 2a correspond to the wavelength of 4 047 Å, those of Fig. 25 are for the wavelength of 2 805 Å, while Fig. 2b corresponds to the wavelength of 2 300 Å. The authors thank Professor N.A. Kaptsov for directing this work. There are 2 figures and 4 references, 2 of which are Russian and 2 English.

ASSOCIATION:

Chair of Electronics of the Physics Faculty of the

Moscow State University im. M.V. Lomonosov

Card 2/3

(Kafedra elektroniki fizicheskogo fakul'teta

Energy Distribution of the Photo-electrons in the Antimony-cuesium Cathode during Various Stages of its Formation

Moskovskogo gosudarstvennogo universiteta im. M.V. Lomonosova)

SUBMITTED: January 19, 1957

AVAILABLE: Library of Congress

Card3/3

9,4160 (1137,1331) 26.1632

83929 S/188/60/000/004/003/014 B005/B060

AUTHORS:

Zheludeva, C. A., Lyuy Bin-i

Energy Distribution of Photoelectrons of an Antimony

Cesium Cathode During Its Treatment With Oxygen TITLE:

Vestnik Moskovskogo universiteta. Seriya 3, fizika,

astronomiya, 1960, No. 4, pp. 26-31 PERIODICAL:

TEXT: The authors of the present paper studied the energy distribution of photoelectrons of an antimony - cesium cathode under three different conditions: prior to sensitization with oxygen, after sensitization up to the peak value of the photocurrent, and after poisoning of the cathode by excess oxygen. This poisoning was carried on until the sensitivity of the cathode became equal to, or a little lower than, that before sensitization. The method of the retarding field in a spherical condenser was used to draw the energy distribution curves of the photoelectrons and to measure the contact potential difference between cathode and anode. The experimental arrangement is described. Fig. 1 has

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Energy Distribution of Photoelectrons of an Antimony - Cesium Cathode During Its Treatment With Oxygen

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three diagrams showing the retardation curves of the antimony - cesium cathode under the three above conditions for the following wavelengths of light: 5461 A, 3130 A, 2804 A, 2482 A. The analysis of these curves yielded the following results: When sensitizing the cathode with oxygen up to a peak value of the photocurrent, the value of the saturation potential rises by 0.6 ± 0.1 v. Investigations revealed that this growth is mainly brought about by a decrease in the thermionic work function of the cathode on a treatment with oxygen. This decrease is about 0.4 ev. On a further treatment of the cathode with oxygen (poisoning), the saturation potential does not change until the sensitivity of the cathode has dropped to the pre-sensitization value. The saturation potential rises little with further poisoning. The mentioned decrease in the thermionic work function of the cathode on sensitization is a consequence of a reduced potential threshold, and fits the change in the threshold frequency of the photoeffect. Fig. 2 shows the spectral characteristics of the antimony - cesium cathode prior to and after sensitization in a (log I, hv) diagram (I = photocurrent; hv = energy in

Card 2/3

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Energy Distribution of Photoelectrons of an Antimony - Cesium Cathode During Its Treatment With Oxygen

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ev). By graphically differentiating the curves shown in Fig. 1, the authors obtained the corresponding curves for the energy distribution of the photoelectrons of the antimony - cesium cathode under the three conditions investigated. Fig. 3 shows these curves for the wavelength 2804 A. It may be seen that the changes in the energy distribution curves of the photoelectrons due to cathode sensitization can be solely explained by a decrease in the work function of the cathode. The energy structure of the cathode need not necessarily change in this connection. It is only on the poisoning of the cathode with excess cxygen and when the cathode loses its sensitivity, that the character of the energy spectrum of the electrons, and thus also the energy structure of the cathode, undergo a change. There are 3 figures and 5 Soviet references.

ASSOCIATION:

Moskovskiy universitet Kafedra elektroniki (k

University, Chair of Electronics)

SUBMITTED:

November 27, 1959

Card 3/3

40405

S/109/62/007/009/015/018 D409/D301

9.4175 AUTHOR:

Zheludeva, G.A.

TITLE:

Influence of fatigue of Sb-Cs cathode on the energy

distribution of photoelectrons

PERIODICAL:

Radiotekhnika i elektronika, v. 7, no. 9, 1962,

1637 - 1642

TEXT: The author studied the influence of cathode fatigue on the energy distribution of photoelectrons by the method of delay curves. The experimental apparatus was similar to that described by the author in an earlier work. The cathodes were prepared on various bases (electrolytically-pure nickel, quartz, molybdenum-glass). The delay curves were plotted at 3 different moments of time: Immediately after the preparation of the experimental tube, one month later, and after 10 months. The study of cathode fatigue started 3-4 weeks after the preparation of the experimental tubes, i.e. after the properties of the cathode were fully stabilized. The fatigue resulted from continuous illumination of the cathode. The measurements of the cathode characteristics were conducted before and Card 1/2

S/109/62/007/009/015/018 D409/D301

Influence of fatigue of Sb-Cs ...

after the fatigue process. The results of the measurements of all the investigated cathodes are similar; the integral sensitivity of the cathodes with nickel base was reduced to half its value as a result of the fatigue (which lasted about 200 hours). In the case of cathodes with molybdenum-glass base, the sensitivity decreased to 68 % of its initial value. The photoelectric work function increased slightly after the fatigue process. The energy spectrum of the photoelectrons after fatigue, differs from the energy spectrum prior to it. The maximum of the energy-distribution curves is shifted towards lower energies, i.e. the energy of the electrons decreases as a result of the fatigue. It is noted that this change in the energy distribution is very similar to that, corresponding to a transition of the cathode from a state of normal activation to a state of de-activation. In both processes (fatigue and de-activation), the relative number of fast electrons increases, whereas the number of slow electrons decreases. The obtained results lead to the conclusion that cathode fatigue is a bulk process, i.e. the fatigue (under the action of light and the applied voltage) involves substantial changes in the cathode material itself. There are 6 figures. SUBMITTED: October 20, 1961 Card 2/2

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1/1553 \$/188/62/000/005/001/008 B102/B108

24,2600

AUTHOR:

Zheludeva, G. A.

TITLE:

Effect of the electron yield depth of a thick antimony-cesium cathode on the nature of their energy distribution

PERIODICAL: Moscow. Universitet. Vestnik. Seriya III. Fizika,
astronomiya, no. 5, 1962, 3 - 9

TEXT: The photoelectron energy distribution under frontal illumination was studied on Sb-Cs cathodes of different thicknesses (test series 1), as well as the effect of the direction of illumination on this energy distribution (test series 2), the cathodes having been produced under completely equal conditions. Cathodes of different thicknesses, however, have different sensitivities. A cathode of ordinary thickness (~1500 Å) is more sensitive than a thinner one, and if the cathode thickness exceeds the effective electron yield depth the "thickness effect" coases to be observable. The sensitivity of a cathode (I) 5300 Å thick was 40 μa/lm and that of a cathode (II) 1200 Å thick was only half as much, although at such a thickness no thickness effect should occur. The measurements were

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Effect of the electron yield ...

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made in the range of wavelengths of 2200 - 4047 A on cathodes of the thicknesses of 490, 1080, 1200, 1600 and 5300 %. All cathodes were produced in the same tube on a quartz backing according to a method described in "Izv. AN SSSR" ser. fiz. 12, no. 2, 126 - 143, 1948. A comparison between the results obtained with the cathodes I and II shows that the photoelectric work function of II is higher than that of I, also & of II was higher by about 0.1 ev than that of I. b is the distance, measured in ev, between the top of the filled band and the electrochemical potential level. energy distributions of the electrons from I and II are different: II emits more numerous fast and fewer slow electrons than I. Hence, it can be concluded that even if the thicknesses of the electrodes under consideration is such as to exclude a thickness effect, they still differ in their most important physical properties. In the second test series the difference in frontal and backward illumination was studied in I and II with light of 2482 R.. Some of the results were unexpected: e.g., the energy spectrum of the "low" electrons contains more numerous fast and fewer slow electrons than that of the "surface" electrons. The energy distributions of the "low" electrons are the same for I and II. In order

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Effect of the electron yield ...

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to explain the results it is assumed that the cathodes have inhomogeneous thicknesses and that the Sb-Cs cathodes act as gray filters on the excited electrons. This inhomogeneity also explains the unexpectedly high photocurrent yield in the case of backward illumination, because the photoactive layer is on the cathode surface and the depth to which the light penetrates is only about 500 %. The "gray filter effect" influences essentially the energy distribution, both with frontal and backward illumination. There are 6 figures.

ASSOCIATION: Kafedra elektroniki (Department. of Electronics)

SUBMITTED: October 28, 1961 (initially), February 10, 1962 (after revision)

Card 3/3